



Solving nonlinear space-time fractional differential equations via ansatz method

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Abstract In this paper, the fractional partial differential equations are defined by modified Riemann-Liouville fractional derivative. With the help of fractional derivative and fractional complex transform, these equations can be converted into the nonlinear ordinary differential equations. By using solitay wave ansatz method, we find exact analytical solutions of the space-time fractional Zakharov Kuznetsov Benjamin Bona Mahony (ZK-BBM) equation, the space-time fractional Klein-Gordon equation and the space-time fractional modified Regularized Long Wave (RLW) equation. This method can be suitable and more powerful for solving other kinds of nonlinear FDEs arising in mathematical physics.

Keywords. Ansatz method, Exact solution, Space-time fractional differential equations.

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1. INTRODUCTION

Fractional differential equations are generalizations of differential equations of integer order. So, in recent years, nonlinear fractional differential equations (FDEs) have gained importance and popularity in various fields of science. These equations appear in a great array of contexts such as in plasma physics, fluid mechanics, nonlinear optics, geochemistry, acoustic waves, hydrodynamics, chemical kinematics, control theory, optical fibers, chemical physics, signal processing, systems identification and many other fields [45, 48, 54].

The study of solitary wave has made remarkable advances in the past decades. Soliton is one of the major areas of research in nonlinear dispersive media. There are two different types of envelope solitons bright and dark. This area of research has made an enormous progress especially in recent years [6, 7, 10, 14–18, 43, 60]. The existence of soliton-type solution for nonlinear PDEs is of particular interest because of their extensive applications in many physics areas. This paper addresses the dynamics

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of soliton propagation through soliton solutions. This leads to a different kind of nonlinear FDEs that describes the dynamics of soliton propagation [29, 30, 49, 50, 52].

There are, in fact, various modern methods of integrability of a variety of nonlinear fractional differential equations. Some of these methods are the exp-function method, the (G'/G) -expansion method, the first integral method, the fractional sub-equation method, the functional variable method, the fractional modified trial equation method, the ansatz method, the modified simple equation method and the modified Kudryashov method [2, 5, 8, 9, 12, 13, 19, 21–28, 31, 32, 34–37, 47, 51, 57, 62].

There are several approaches to the generalization of the notion of differentiation to fractional orders e.g. Grünwald–Letnikov, Caputo and Riemann–Liouville [20, 58]. Modified Riemann–Liouville derivative is defined a local fractional derivative by Jumarie [40]. The definition and some properties for the Jumarie’s derivative of order α are listed as follows [41]

$$D_w^\alpha f(w) = \begin{cases} \frac{1}{\Gamma(1-\alpha)} \frac{d}{dw} \int_0^w \frac{f(\tau)-f(0)}{(w-\tau)^\alpha} d\tau, & 0 < \alpha < 1, \\ (f^{(n)}(w))^{(\alpha-n)}, & n \leq \alpha \leq n+1, n \geq 1. \end{cases} \quad (1.1)$$

$$D_w^\alpha w^\gamma = \frac{\Gamma(1+\gamma)}{\Gamma(1+\gamma-\alpha)} w^{\gamma-\alpha}, \quad \gamma > 0, \quad (1.2)$$

$$D_w^\alpha (c) = 0, \quad c = \text{constant}, \quad (1.3)$$

$$D_w^\alpha \{af(w) + bg(w)\} = aD_w^\alpha f(w) + bD_w^\alpha g(w), \quad (1.4)$$

where $a \neq 0$ and $b \neq 0$ are constants. Now, we will give main steps of methodology.

Step 1: We consider the following general nonlinear space-time FDE of the type

$$H(u, D_t^\alpha u, D_x^\alpha u, D_t^{2\alpha} u, D_t^\alpha D_x^\alpha u, D_x^{2\alpha} u, \dots) = 0, \quad 0 < \alpha < 1, \quad (1.5)$$

where u is an unknown function, and H is a polynomial of u and its partial fractional derivatives.

Step 2: By using fractional complex transform

$$\begin{aligned} u(x, t) &= f(\tau), \\ \tau &= \frac{kx^\alpha}{\Gamma(1+\alpha)} - \frac{ct^\alpha}{\Gamma(1+\alpha)}, \end{aligned} \quad (1.6)$$

where $k \neq 0$ and $c \neq 0$ are constants and by using the chain rule

$$\begin{aligned} D_t^\alpha u &= \sigma_t \frac{\partial f}{\partial \tau} D_t^\alpha \tau, \\ D_x^\alpha u &= \sigma_x \frac{\partial f}{\partial \tau} D_x^\alpha \tau, \end{aligned} \quad (1.7)$$

where σ_t, σ_x are called the sigma indexes [33] and it can take $\sigma_t = \sigma_x = L$, where L is a constant.

Step 3: When we substitute (1.6) with (1.2) and (1.7) into (1.5), we get following nonlinear ODE,

$$N\left(U, \frac{df}{d\tau}, \frac{d^2 f}{d\tau^2}, \frac{d^3 f}{d\tau^3}, \dots\right) = 0. \quad (1.8)$$



2. APPLICATIONS

2.1. **The space-time fractional ZK-BBM equation.** Let us consider, the space-time fractional ZK-BBM equation [3]

$$D_t^\alpha u + D_x^\alpha u - 2auD_x^\alpha u - bD_t^\alpha(D_x^{2\alpha}u) = 0, \tag{2.1}$$

where a and b are arbitrary constants. It arises as a description of gravity water waves in the long-wave regime. Alzaidy solved this equation by a fractional sub-equation method in [3] and obtained three types of exact analytical solutions. Bekir et al. have applied the functional variable method to obtain new periodic and hyperbolic solutions of Eq.(2.1) in [11]. We will use the ansatz method to obtain the exact solutions with the help of ansatz method. In order to solve Eq.(2.1), we use the transformation (1.6) then integrating Eq.(2.1) with respect to τ and setting the integration constant equal to zero, we have

$$(k - c)f - akf^2 + bck^2L^2f'' = 0. \tag{2.2}$$

To obtain bright soliton solution of Eq.(2.2),

$$f(\tau) = A \operatorname{sech}^p \tau, \tag{2.3}$$

where

$$\tau = \frac{kx^\alpha}{\Gamma(1 + \alpha)} - \frac{ct^\alpha}{\Gamma(1 + \alpha)}, \tag{2.4}$$

which k , c and A are constant coefficients. From the ansatz (2.3) and (2.4), we get

$$\frac{d^2f(\tau)}{d\tau^2} = Ap^2 \operatorname{sech}^p \tau - Ap(p + 1) \operatorname{sech}^{p+2} \tau, \tag{2.5}$$

and

$$f^2(\tau) = A^2 \operatorname{sech}^{2p} \tau. \tag{2.6}$$

Thus, substituting the ansatz (2.3)-(2.6) into Eq.(2.2), yields to

$$\begin{aligned} &(k - c)A \operatorname{sech}^p \tau - akA^2 \operatorname{sech}^{2p} \tau \\ &- bck^2L^2Ap^2 \operatorname{sech}^p \tau - bck^2L^2Ap(p + 1) \operatorname{sech}^{p+2} \tau \\ &= 0. \end{aligned} \tag{2.7}$$

From (2.7), when we equate exponents $p + 2$ and $2p$, that leads to $p = 2$. From (2.7), setting the coefficients of $\operatorname{sech}^{p+2} \tau$ and $\operatorname{sech}^{2p} \tau$ terms to zero,

$$-akA^2 - bck^2L^2Ap(p + 1) = 0, \tag{2.8}$$

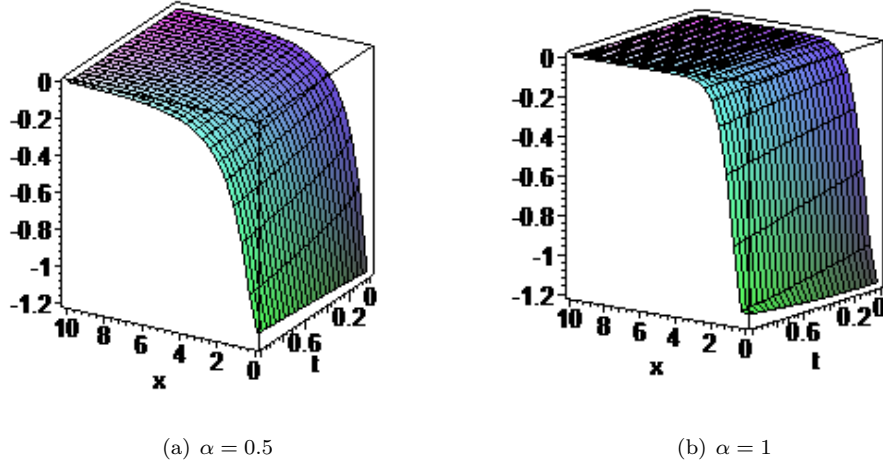
then we obtain

$$A = -\frac{bckL^2p(p + 1)}{a}. \tag{2.9}$$

We find, from setting the coefficients of $\operatorname{sech}^p \tau$ terms in Eq.(2.7) to zero

$$(k - c)A - bck^2L^2Ap^2 = 0, \tag{2.10}$$



FIGURE 1. Shape of solution for (2.12) with $k = 1, a = 1, b = 1, L = 1$.

also we get

$$c = \frac{k}{bk^2L^2p^2 + 1}. \quad (2.11)$$

Finally; when we use $p = 2$, we get the bright soliton solution for the space-time fractional ZK-BBM equation as follow:

$$u(x, t) = -\frac{6bckL^2A}{a} \operatorname{sech}^2 \left(\frac{kx^\alpha}{\Gamma(1 + \alpha)} - \frac{kt^\alpha}{(4bk^2L^2 + 1)\Gamma(1 + \alpha)} \right). \quad (2.12)$$

The solution.(2.12) is represented in Figure 1, within the interval $0 < x < 10$ and $0 < t < 1$.

Remark 1. The solution (2.12) is not given in [3, 11] and have not been reported by other authors in the literature.

2.2. The nonlinear space-time fractional Klein-Gordon equation. We next consider the nonlinear space-time fractional Klein-Gordon (KG) equation [39]

$$D_{tt}^{2\alpha} u - D_{xx}^{2\alpha} u + \gamma u - \beta u^2 = 0. \quad (2.13)$$

where α, β are nonzero constant. This equation describes many types of nonlinearities and plays a significant role in several real world applications such as the solid state physics, nonlinear optics and quantum field theory. Baleanu and his colleagues have found many new types of exact travelling wave solutions of KG equation by using the auxiliary equation method by using sub-equation method, and obtained new exact solutions of equation (2.13) containing hyperbolic, trigonometric and rational functions. In [11], the functional variable method successfully applied to finding periodic and hyperbolic solutions of the fractional KG equation by Bekir at al. When



$\alpha = 1$, equation (2.13) is called the quadratic nonlinear Klein–Gordon equation and there are a lot of studies for this equation [4, 38, 46, 53, 59, 61, 63].

Now we consider the nonlinear fractional KG equation. With the same process as in the previous example, we obtain following ODE

$$(c^2L^2 - k^2L^2)f'' + \gamma U - \beta U^2 = 0, \tag{2.14}$$

where $f' = \frac{df}{d\tau}$.

From the ansatz (2.3) and (2.4), we obtain necessary derivatives. Then, substituting them into Eq.(2.14), yields to

$$\begin{aligned} &(c^2L^2 - k^2L^2)Ap^2 \operatorname{sech}^p \tau - (c^2L^2 - k^2L^2)Ap(p + 1) \operatorname{sech}^{p+2} \tau \\ &+ \gamma A \operatorname{sech}^p \tau - \beta A^2 \operatorname{sech}^{2p} \tau \\ &= 0. \end{aligned} \tag{2.15}$$

From (2.15), equating exponents $p + 2$ and $2p$, that gives $p = 2$. From (2.15), setting the coefficients of $\operatorname{sech}^{p+2} \tau$ and $\operatorname{sech}^{2p} \tau$ terms to zero,

$$(c^2L^2 - k^2L^2)Ap(p + 1) + \beta A^2 = 0, \tag{2.16}$$

by use (4.10), we obtain

$$A = -\frac{(c^2L^2 - k^2L^2)p(p + 1)}{\beta}. \tag{2.17}$$

Analogously, from setting the coefficients of $\operatorname{sech}^p \tau$ terms in Eq.(2.15) to zero, we have

$$(c^2L^2 - k^2L^2)Ap^2 + \gamma A = 0, \tag{2.18}$$

then we get

$$c = \mp \frac{\sqrt{p^2k^2L^2 + \gamma}}{pL}. \tag{2.19}$$

Consequently, we can determine the bright soliton solution of (2.13) as with $p = 2$,

$$u(x, t) = -\frac{6(c^2L^2 - k^2L^2)A}{\beta} \operatorname{sech}^2 \left(\frac{kx^\alpha}{\Gamma(1 + \alpha)} \pm \frac{\sqrt{4k^2L^2 + \gamma}t^\alpha}{2L\Gamma(1 + \alpha)} \right). \tag{2.20}$$

Also, Eq.(2.19) implies the domain restrictions $4k^2L^2 + \gamma > 0$.

We plot the solutions of Eq.(2.20) for this equation in Figure 2 within the interval $0 < x < 100$ and $0 < t < 100$.

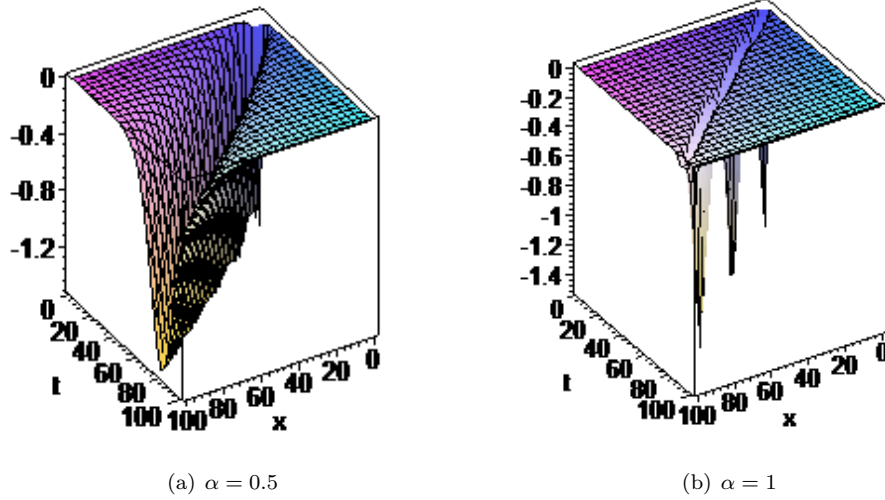
Remark 2. Comparing our result with Baleanu’s and Bekir’s [11, 39] results, it can be seen that the result is new.

2.3. The space-time fractional modified RLW equation. This equation has the form [42]

$$D_t^\alpha u + vD_x^\alpha u + \mu u^2 D_x^\alpha u - \varepsilon D_t^\alpha D_x^{2\alpha} u = 0, \tag{2.21}$$

where α describing the order of the fractional derivatives $0 < \alpha \leq 1$ and v, μ and ε are all constants. Kaplan et al. solved this equation by the modified simple equation



FIGURE 2. Shape of solution for (2.20) with $k = 1, L = 1, \beta = 1, \gamma = 1$.

method [42]. Abdel-Salam and Gumma have obtained abundant types of exact analytical solutions including generalized trigonometric and hyperbolic functions solutions of this equation with the improved fractional Riccati expansion method in [1]. The modified RLW equation is considered as an alternative to the modified KdV equation. This equation is modeled to govern a large number of physical phenomena such as transverse waves in shallow water and magneto hydrodynamic waves in plasma and phonon packets in nonlinear crystals [44, 55, 56].

When we substitute (1.6) with (1.2) and (1.7) into (2.21), integrating Eq.(2.1) with respect to τ and setting the integration constant equal to zero, Eq.(2.21) can reduced into an ODE

$$(kv - c)f + \frac{\mu k}{3}f^3 + \varepsilon ck^2 L^2 f'' = 0, \quad (2.22)$$

where $f' = \frac{df}{d\tau}$.

From the ansatz (2.3) and (2.4), we obtain necessary derivatives. Then, substituting them into Eq.(2.22), yields to

$$\begin{aligned} & (kv - c)A \operatorname{sech}^p \tau + \frac{\mu k}{3}A^3 \operatorname{sech}^{3p} \tau \\ & + \varepsilon ck^2 L^2 A p^2 \operatorname{sech}^p \tau - \varepsilon ck^2 L^2 A p(p + 1) \operatorname{sech}^{p+2} \tau \\ & = 0. \end{aligned} \quad (2.23)$$

From (2.23), if we equate the exponents $p + 2$ and $3p$, we have

$$p = 1. \quad (2.24)$$



When we set, coefficients of $\operatorname{sech}^{p+2} \tau$ and $\operatorname{sech}^{3p} \tau$ terms to zero in Eq.(2.23), we get

$$\frac{\mu k}{3} A^3 - \varepsilon c k^2 L^2 A p(p+1) = 0, \quad (2.25)$$

by use (2.24) and after some calculations, we have

$$A = \mp L \sqrt{\frac{6\varepsilon k c}{\mu}}, \quad \mu \neq 0. \quad (2.26)$$

Again from setting coefficients of $\operatorname{sech}^p \tau$ terms in Eq.(2.23) to zero

$$(k v - c) A + \varepsilon c k^2 L^2 A p^2 = 0, \quad (2.27)$$

we obtain

$$c = \frac{v k}{1 - \varepsilon k^2 L^2}. \quad (2.28)$$

From (2.28) it is important to note that

$$4\varepsilon k^2 L^2 \neq 1. \quad (2.29)$$

Thus finally, we can determine the bright soliton solution of (2.21) as with $p = 1$,

$$u(x, t) = A \operatorname{sech} \left(\frac{k x^\alpha}{\Gamma(1 + \alpha)} - \frac{c t^\alpha}{\Gamma(1 + \alpha)} \right), \quad (2.30)$$

where the A is given by (2.26) and the c is given by (2.28). We plot the solutions of Eq.(2.30) for this equation in Figure 3 within the interval $0 < x < 100$ and $0 < t < 100$.

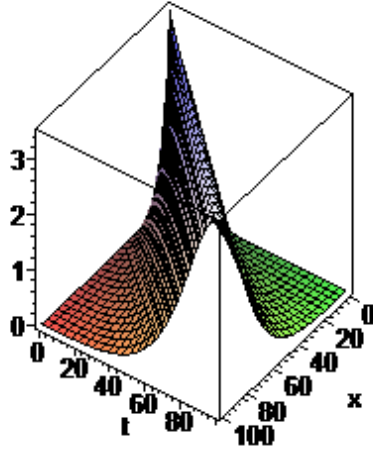
Remark 3. Note that solution (2.30) is quite different from the travelling wave solutions found in [1, 42].

3. CONCLUSION

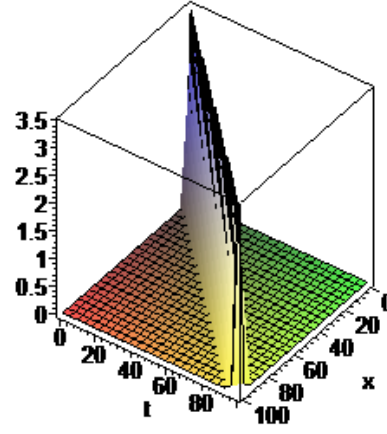
In this paper, the ansatz method is used to obtain the bright soliton solution of the nonlinear FDEs with Jumarie's modified Riemann–Liouville derivative. In general, there exist no method that yields soliton solutions for fractional differential equations. But, a fractional complex transform is adopted in this paper to convert such equations into classical partial differential equations. We succeeded in extracting soliton solutions for the space-time fractional ZK-BBM equation, the space-time fractional Klein-Gordon equation and the space-time fractional modified RLW equation. As a result, some new exact solutions for them have been successfully found. Being concise and powerful, the proposed method can be applied to solve other nonlinear FDEs and systems. All the solutions reported above have been verified using the symbolic computation system Maple.



FIGURE 3. Shape of solution for (2.30) with $k = 1, \tau = 2, v = -1, L = 1, \mu = 1$.



(a) $\alpha = 0.5$



(b) $\alpha = 1$

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