# A compact difference scheme for time-fractional Black-Scholes equation with time-dependent parameters under the CEV model: American options 

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#### Abstract

The Black-Scholes equation is one of the most important mathematical models in option pricing theory, but this model is far from market realities and cannot show memory effect in the financial market. This paper investigates an American option based on a time-fractional Black-Scholes equation under the constant elasticity of variance (CEV) model, which parameters of interest rate and dividend yield supposed as deterministic functions of time, and the price change of the underlying asset follows a fractal transmission system. This model does not have a closed-form solution; hence, we numerically price the American option by using a compact difference scheme. Also, we compare the time-fractional Black-Scholes equation under the CEV model with its generalized Black-Scholes model as $\alpha=1$ and $\beta=0$. Moreover, we demonstrate that the introduced difference scheme is unconditionally stable and convergent using Fourier analysis. The numerical examples illustrate the efficiency and accuracy of the introduced difference scheme.


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## 1. Introduction

According to exercise time, we have two types of options: American and European options. An American option can be exercised at any time before the maturity date, but a European option can only be exercised at the maturity date. The American option must be worth at least as much as its European equivalent and it can never be worth less than its payoff. This option is more valuable than the European option, because investors have the freedom to exercise their option at any time during the contract while holder of European option can only exercise at maturity.

The pricing problem of American options is a very common class of optimalstopping problems that leads to a free boundary value problem. In 1965, McKean [31] was the first person who worked on the American options and derived a free-boundary problem to determine the price of American options. In 1976, Moerbeke [36] further studied on the properties of the optimal exercise boundary and extended its analysis. Then, numerical methods for solving the free boundary problem are developed by Brennan and Schwartz (in 1977) [4], Schwartz (in 1977) [42] and Courtadon (in 1982) [9].

Analytical solutions of European options can easily be found, while there exists no exact solution for the American options, because its price depends on the history of the underlying asset price and its present value. There are various numerical methods to obtain an optimal solution of the free boundary problem, that one of the most common methods is binomial methods. Also, Muthuraman [38], Chockalingam, and Muthuraman [8] used the moving boundary approach to price American option. Chen et al. [6] introduced a predictor-corrector approach based on the spectralcollocation method to price the American options under the finite moment log-stable model. Moradipour and Yousefi used collocation methods to solve the Black-Scholes equation for American option pricing [37].

The classical Black-Scholes model is under some restrictive assumptions that it makes to weaken. Hence, some developed models are required, such as the fractional Black-Scholes model [17, 19, 32, 51], the Black-Scholes model with the transaction costs [27, 44, 47], the stochastic volatility model [18, 30], and the jump-diffusion model [22, 35]. The classical Black-Scholes models cannot show memory effect in financial systems as well. In the last decade, the fractional Black-Scholes models are used to describe the effect of trend and noise memory in financial pricing. Li et al. [25] priced the European option based on the fractional order stochastic differential equation model and derived the trend memory in stock price process when the Hurst index is between 0.5 and 1. Moreover, they presented a new approach in option pricing that it leads to a better result than the classic model and stochastic model with fractional Brownian motion when the stock prices are simulated by Monte Carlo simulation. Liu and Chang [27] presented a formula for European option pricing with transaction costs based on the fractional Black-Scholes model and showed that the price of the European option decreases with the increase of the Hurst index. Zhang et al. [52] investigated the tempered fractional Black-Scholes equation with the numerical simulation for pricing of a European double barrier option under three models of Finite Moment Log Stable, KoBoL, and CGMY. Mehrdoust et al. [34] considered
the mixed fractional Heston model to show long-range dependence and exhibited that the Euler discretization method on this model has strong convergence. Also, they estimated the American put option price under this model. Besides, other classes of fractional Black-Scholes equations are introduced by Jumarie [20] and Farhadi et al. [15].

In this paper, we consider the time-fractional Black-Scholes equation to price the American call options whose parameters: interest rate, dividend yield, and volatility are as functions, while they are considered constant in the classical models, and these assumptions of the problem are closer to the actual model of the market. In many markets, stock price volatility is increasing when the stock price is decreasing. For modeling this phenomenon, we cannot use the classical Black-Scholes model and will have to use the generalized models. Including these models: local volatility models [16] and stochastic volatility models [43]. One of the most famous of these models is the CEV model that was introduced for the first time by Cox [10] in 1975 to record the inverse relationship between stock prices and their volatility. The volatility of this model without introducing any additional random process is a function of the stock price and two parameters of $\beta$ and $\delta$ that are called elasticity of volatility and scale parameter fixing the initial instantaneous volatility, respectively. An important parameter of this model is the elasticity of volatility that controls the relationship between volatility and asset price. This model has been widely used in many areas, including: determining the value of American options [1, 48, 49, 53], Asian options [23, 34], Barrier options [28, 46], and Lookback options [3, 13, 21]. Since the closedform solution does not exist to price the American option, our main aim of this paper is to obtain an American option price under the CEV model with a compact difference scheme. In the following, we investigate the stability and convergence of the introduced scheme. Staelen and Hendy [14] have already used this scheme to price the double barrier options for time-fractional Black-Scholes model whose parameters are constant.

The main body of this paper is organized as follows: In section 2, we formulate time-fractional Black-Scholes equation under the CEV model for pricing American option. In section 3, we construct the compact difference scheme to price the American option numerically. In section 4, solvability, unconditionally stability and convergence of introduced difference scheme are illustrated using Fourier analysis. In section 5, numerical examples demonstrate the efficiency of compact difference scheme to solve time-fractional Black-Scholes equation with time-dependent parameters. Finally, the paper ends with remarks and conclusions.

## 2. American option pricing model

In this paper, the underlying asset price is considered under CEV model as follows (see [11])

$$
\begin{equation*}
d S_{t}=(r(t)-D(t)) S_{t} d t+\delta S_{t}^{\beta+1} d W_{t}, \tag{2.1}
\end{equation*}
$$

where functions $r(t)$ and $D(t)$ are the risk-free interest rate and dividend yield, respectively. $W_{t}$ is standard Brownian motion, $d S_{t}$ is the change in the stock price $S$ over the short increment of time $d t$, also, $\beta$ and $\delta\left(\delta=\sigma_{0} S_{0}^{-\beta}\right)$ are positive constants.


Noticing that in the case of $\beta>0$, the local volatility function, $\sigma(S)=\delta S^{\beta}$, does not remain bounded as $S \rightarrow+\infty$. Therefore, we consider $\beta<0$ that $\sigma(S)$ remains bounded and decreases as the asset price increases. In particular, when $\beta=0$ the volatility $\sigma(S)=\delta S^{\beta}$ is constant and the CEV model (2.1) turns into the lognormal diffusion model which is the generalized Black-Scholes model [2]. Otherwise, the following partial differential equation will be obtained using the Itô Lemma

$$
\frac{\partial C(S, t)}{\partial t}+\frac{1}{2} \delta^{2} S^{2 \beta+2} \frac{\partial^{2} C(S, t)}{\partial S^{2}}+(r(t)-D(t)) S \frac{\partial C(S, t)}{\partial S}-r(t) C(S, t)=0
$$

which $C(S, t)$ is the value of the American call option with asset price $S_{t}$ in the moment $t$. Under this assumption that the price change of the underlying follows a fractal transmission system, we can replace the time derivative $\frac{\partial C}{\partial t}$ with the fractional derivative $\frac{\partial^{\alpha} C}{\partial t^{\alpha}}$ (see, e.g., $[26,7]$ ) as follows

$$
\begin{equation*}
\frac{\partial^{\alpha} C(S, t)}{\partial t^{\alpha}}+\frac{1}{2} \delta^{2} S^{2 \beta+2} \frac{\partial^{2} C(S, t)}{\partial S^{2}}+(r(t)-D(t)) S \frac{\partial C(S, t)}{\partial S}-r(t) C(S, t)=0 \tag{2.2}
\end{equation*}
$$

where $\frac{\partial^{\alpha} C}{\partial t^{\alpha}}$ is the modified right Riemann-Liouville fractional derivative and is defined as

$$
\frac{\partial^{\alpha} C(S, t)}{\partial t^{\alpha}}=\left\{\begin{array}{lr}
\frac{1}{\Gamma(1-\alpha)} \frac{\partial}{\partial t} \int_{t}^{T} \frac{C(S, \xi)-C(S, T)}{(\xi-t)^{\alpha}} d \xi, & 0<\alpha<1 \\
\frac{\partial C(S, t)}{\partial t}, & \alpha=1
\end{array}\right.
$$

With assumption, diffusion of the option price depends on the history of the time to maturity. By using change variable $\tau=T-t$ and $V(S, \tau) \equiv C(S, T-\tau)=C(S, t)$, we turn backward problem to forward problem and show the relationship $\frac{\partial^{\alpha} C}{\partial t^{\alpha}}$ with the Caputo fractional derivative. We calculate as follow

$$
\begin{align*}
& \frac{\partial^{\alpha} C(S, t)}{\partial t^{\alpha}}=\frac{1}{\Gamma(1-\alpha)} \frac{\partial}{\partial t} \int_{t}^{T} \frac{C(S, \rho)-C(S, T)}{(\rho-t)^{\alpha}} d \rho \\
& \quad \stackrel{\rho=T-\zeta}{=} \frac{1}{\Gamma(1-\alpha)} \frac{\partial}{\partial \tau} \int_{\tau}^{0} \frac{C(S, T-\zeta)-C(S, T)}{(\tau-\zeta)^{\alpha}} d \zeta \\
& \quad=\frac{-1}{\Gamma(1-\alpha)} \frac{\partial}{\partial \tau} \int_{0}^{\tau} \frac{V(S, \zeta)-V(S, 0)}{(\tau-\zeta)^{\alpha}} d \zeta \\
& =\frac{-1}{\Gamma(1-\alpha)} \frac{\partial}{\partial \tau} \int_{0}^{\tau}[V(S, \zeta)-V(S, 0)] d \frac{(\tau-\zeta)^{1-\alpha}}{-(1-\alpha)} \\
& =\frac{-1}{\Gamma(1-\alpha)} \frac{\partial}{\partial \tau} \int_{0}^{\tau} \frac{(\tau-\zeta)^{1-\alpha}}{1-\alpha} \frac{\partial V(S, \zeta)}{\partial \zeta} d \zeta  \tag{2.3}\\
& =\frac{-1}{\Gamma(1-\alpha)} \int_{0}^{\tau} \frac{1}{(\tau-\zeta)^{\alpha}} \frac{\partial V(S, \zeta)}{\partial \zeta} d \zeta .
\end{align*}
$$

Furthermore, right side of (2.3) is definition of the left-hand side Caputo fractional as follows [40]

$$
{ }_{0}^{C} D_{\tau}^{\alpha} V(S, \tau):=\frac{1}{\Gamma(1-\alpha)} \int_{0}^{\tau} \frac{1}{(\tau-\zeta)^{\alpha}} \frac{\partial V(S, \zeta)}{\partial \zeta} d \zeta .
$$

Thus,

$$
\begin{equation*}
\frac{\partial^{\alpha} C(S, t)}{\partial t^{\alpha}}=-{ }_{0}^{C} D_{\tau}^{\alpha} V(S, \tau) \tag{2.4}
\end{equation*}
$$

Due to the (2.2), (2.4), $\frac{\partial C}{\partial S}=\frac{\partial V}{\partial S}$ and $\frac{\partial^{2} C}{\partial S^{2}}=\frac{\partial^{2} V}{\partial S^{2}}$ the estimation of American call option can be formulated as a time-fractional free boundary problem:

$$
\begin{align*}
& { }_{0}^{C} D_{\tau}^{\alpha} V(S, \tau)=\frac{1}{2} \delta^{2} S^{2 \beta+2} \frac{\partial^{2} V}{\partial S^{2}}+(r(\tau)-D(\tau)) S \frac{\partial V}{\partial S}-r(\tau) V,  \tag{2.5}\\
& \quad 0<\tau<T, \quad 0<S<S_{f}(\tau), \\
& V(S, 0)=\max (S-K, 0),  \tag{2.6}\\
& V(0, \tau)=0,  \tag{2.7}\\
& V\left(S_{f}(\tau), \tau\right)=S_{f}(\tau)-K,  \tag{2.8}\\
& \frac{\partial V\left(S_{f}(\tau), \tau\right)}{\partial S}=1, \tag{2.9}
\end{align*}
$$

where $K$ is "exercise price" or "strike price" and $T$ is "expiration date" or "maturity date".

## 3. Compact difference scheme

In this section, we derive a numerical solution for problem (2.5) by using compact difference scheme. At first, we define following uniform time and space mesh for any positive integers $M$ and $N$ ( $S_{\max }$ is a sufficiently large number)

$$
\begin{aligned}
& \Delta \tau=\frac{T}{M}, \quad \tau_{n}=n \Delta \tau, \quad n=0,1, \ldots, M, \\
& \Delta S=\frac{S_{\max }}{N}, \quad S_{j}=j \Delta S, \quad j=0,1, \ldots, N .
\end{aligned}
$$

By using above uniform mesh, we can formulate time-fractional derivative ${ }_{0}^{C} D_{\tau}^{\alpha} V(S, \tau)$ of problem (2.5) at point $\left(S_{j}, \tau_{n}\right)$ with notations $V_{j}^{n}=V\left(S_{j}, \tau_{n}\right)(j=0,1, \ldots, N ; n=$ $0,1 \ldots, M)$ as

$$
\begin{equation*}
{ }_{0}^{C} D_{\tau}^{\alpha} V\left(S_{j}, \tau_{n}\right)=\varphi \sum_{k=1}^{n} \psi_{k}\left(V_{j}^{n-k+1}-V_{j}^{n-k}\right)+O\left(\Delta \tau^{2-\alpha}\right), \tag{3.1}
\end{equation*}
$$

with

$$
\varphi=\frac{1}{\Gamma(2-\alpha)} \frac{1}{\Delta \tau^{\alpha}}, \quad \psi_{k}=k^{1-\alpha}-(k-1)^{1-\alpha}, \quad k=1, \ldots, n
$$

and approximate spatial derivatives based on Taylor expansion of $V \in \mathbb{C}^{4}(0,+\infty)$ as follows

$$
\begin{align*}
& \frac{\partial V\left(S_{j}, \tau_{n}\right)}{\partial S}=\underbrace{\frac{V\left(S_{j+1}, \tau_{n}\right)-V\left(S_{j-1}, \tau_{n}\right)}{2 \Delta S}}_{:=\delta_{S} V\left(S_{j}, \tau_{n}\right)}-\frac{\Delta S^{2}}{6} \frac{\partial^{3} V\left(S_{j}, \tau_{n}\right)}{\partial S^{3}}  \tag{3.2a}\\
& \quad+O\left(\Delta S^{4}\right),
\end{align*}
$$

$$
\begin{align*}
& \frac{\partial^{2} V\left(S_{j}, \tau_{n}\right)}{\partial S^{2}}=\underbrace{\frac{V\left(S_{j-1}, \tau_{n}\right)-2 V\left(S_{j}, \tau_{n}\right)+V\left(S_{j+1}, \tau_{n}\right)}{\Delta S^{2}}}_{:=\delta_{S}^{2} V\left(S_{j}, \tau_{n}\right)}  \tag{3.2b}\\
& \quad-\frac{\Delta S^{2}}{12} \frac{\partial^{4} V\left(S_{j}, \tau_{n}\right)}{\partial S^{4}}+O\left(\Delta S^{4}\right) .
\end{align*}
$$

By substituting (3.2a) and (3.2b) in (2.5) at grid point $\left(S_{j}, \tau_{n}\right)$, we obtain

$$
\begin{align*}
& a S_{j}^{2 \beta+2} \delta_{S}^{2} V\left(S_{j}, \tau_{n}\right)+b_{n} S_{j} \delta_{S} V\left(S_{j}, \tau_{n}\right)-P_{j}^{n}=g\left(S_{j}, \tau_{n}\right)  \tag{3.3}\\
& P_{j}^{n}=\frac{\Delta S^{2}}{12}\left(a S_{j}^{2 \beta+2} \frac{\partial^{4} V\left(S_{j}, \tau_{n}\right)}{\partial S^{4}}+2 b_{n} S_{j} \frac{\partial^{3} V\left(S_{j}, \tau_{n}\right)}{\partial S^{3}}\right)+O\left(\Delta S^{4}\right) \\
& g\left(S_{j}, \tau_{n}\right)={ }_{0}^{C} D_{\tau}^{\alpha} V\left(S_{j}, \tau_{n}\right)+c_{n} V\left(S_{j}, \tau_{n}\right) \\
& a=\frac{\delta^{2}}{2}, \quad b(\tau)=r(T-\tau)-D(T-\tau), \quad c(\tau)=r(T-\tau)
\end{align*}
$$

Also, from (2.5), (3.2a) and (3.2b), we have

$$
\begin{align*}
& \frac{\partial^{3} V\left(S_{j}, \tau_{n}\right)}{\partial S^{3}}=-\frac{2 \beta+2}{a} S_{j}^{-2 \beta-3}\left[g\left(S_{j}, \tau_{n}\right)-b_{n} S_{j} \delta_{S} V\left(S_{j}, \tau_{n}\right)\right] \\
& \quad+\frac{1}{a} S_{j}^{-2 \beta-2}\left[\delta_{S} g\left(S_{j}, \tau_{n}\right)-b_{n} \delta_{S} V\left(S_{j}, \tau_{n}\right)-b_{n} S_{j} \delta_{S}^{2} V\left(S_{j}, \tau_{n}\right)\right]  \tag{3.4}\\
& \quad+O\left(\Delta S^{2}\right)
\end{align*}
$$

and

$$
\begin{align*}
& \frac{\partial^{4} V\left(S_{j}, \tau_{n}\right)}{\partial S^{4}}=\frac{2 \beta+2}{a} S_{j}^{-2 \beta-4}\left[2 \beta+3+\frac{b_{n}}{a} S_{j}^{-2 \beta}\right] \\
& \quad \times\left[g\left(S_{j}, \tau_{n}\right)-b_{n} S_{j} \delta_{S} V\left(S_{j}, \tau_{n}\right)\right]-\frac{1}{a} S_{j}^{-2 \beta-3}\left[4(\beta+1)+\frac{b_{n}}{a} S_{j}^{-2 \beta}\right]  \tag{3.5}\\
& \quad \times\left[\delta_{S} g\left(S_{j}, \tau_{n}\right)-b_{n} \delta_{S} V\left(S_{j}, \tau_{n}\right)-b_{n} S_{j} \delta_{S}^{2} V\left(S_{j}, \tau_{n}\right)\right] \\
& \quad-\frac{1}{a} S_{j}^{-2 \beta-2}\left[\delta_{S}^{2} g\left(S_{j}, \tau_{n}\right)-2 b_{n} \delta_{S}^{2} V\left(S_{j}, \tau_{n}\right)\right]+O\left(\Delta S^{2}\right)
\end{align*}
$$

Now, we substitute (3.4) and (3.5) in $P_{j}^{n}$

$$
\begin{aligned}
P_{j}^{n} & =\frac{(\beta+1)}{6 S_{j}^{2}} \Delta S^{2}\left[2 \beta+3-\frac{b_{n}}{a} S_{j}^{-2 \beta}\right] g\left(S_{j}, \tau_{n}\right) \\
& +\frac{\Delta S^{2}}{12}\left[-\frac{4(\beta+1)}{S_{j}}+\frac{b_{n}}{a} S_{j}^{-2 \beta-1}\right] \delta_{S} g\left(S_{j}, \tau_{n}\right)+\frac{\Delta S^{2}}{12} \delta_{S}^{2} g\left(S_{j}, \tau_{n}\right) \\
& +\frac{b_{n}(2 \beta+1)}{12 S_{j}} \Delta S^{2}\left[-2 \beta-2+\frac{b_{n}}{a} S_{j}^{-2 \beta}\right] \delta_{S} V\left(S_{j}, \tau_{n}\right) \\
& +\frac{b_{n}}{12} \Delta S^{2}\left[4 \beta+2-\frac{b_{n}}{a} S_{j}^{-2 \beta}\right] \delta_{S}^{2} V\left(S_{j}, \tau_{n}\right)+O\left(\Delta S^{4}\right)
\end{aligned}
$$

Then, we set $P_{j}^{n}$ in Eq. (3.3) and after simplification we have

$$
\begin{aligned}
& {\left[a S_{j}^{2 \beta+2}-\frac{b_{n}}{12} \Delta S^{2}\left(4 \beta+2-\frac{b_{n}}{a} S_{j}^{-2 \beta}\right)\right] \delta_{S}^{2} V\left(S_{j}, \tau_{n}\right)} \\
& \quad+\left[b_{n} S_{j}+\frac{b_{n}(2 \beta+1)}{12 S_{j}} \Delta S^{2}\left(2 \beta+2-\frac{b_{n}}{a} S_{j}^{-2 \beta}\right)\right] \delta_{S} V\left(S_{j}, \tau_{n}\right) \\
& =\left[1+\frac{(\beta+1)}{6 S_{j}^{2}} \Delta S^{2}\left(2 \beta+3-\frac{b_{n}}{a} S_{j}^{-2 \beta}\right)\right] g\left(S_{j}, \tau_{n}\right) \\
& \quad+\frac{\Delta S^{2}}{12}\left[-\frac{4(\beta+1)}{S_{j}}+\frac{b_{n}}{a} S_{j}^{-2 \beta-1}\right] \delta_{S} g\left(S_{j}, \tau_{n}\right) \\
& \quad+\frac{\Delta S^{2}}{12} \delta_{S}^{2} g\left(S_{j}, \tau_{n}\right)+O\left(\Delta S^{4}\right)
\end{aligned}
$$

We define the coefficients of the above equation as follows

$$
\begin{align*}
& \gamma_{j}^{n} \delta_{S}^{2} V\left(S_{j}, \tau_{n}\right)+\xi_{j}^{n} \delta_{S} V\left(S_{j}, \tau_{n}\right)=\theta_{j}^{n} g\left(S_{j}, \tau_{n}\right)+w_{j}^{n} \delta_{S} g\left(S_{j}, \tau_{n}\right) \\
&+\frac{\Delta S^{2}}{12} \delta_{S}^{2} g\left(S_{j}, \tau_{n}\right)+O\left(\Delta S^{4}\right)  \tag{3.6}\\
& \gamma_{j}^{n}:=a S_{j}^{2 \beta+2}-\frac{b_{n}}{12} \Delta S^{2}\left(4 \beta+2-\frac{b_{n}}{a} S_{j}^{-2 \beta}\right) \\
& \xi_{j}^{n}:=b_{n} S_{j}+\frac{b_{n}(2 \beta+1)}{12 S_{j}} \Delta S^{2}\left(2 \beta+2-\frac{b_{n}}{a} S_{j}^{-2 \beta}\right) \\
& \theta_{j}^{n}:=1+\frac{(\beta+1)}{6 S_{j}^{2}} \Delta S^{2}\left(2 \beta+3-\frac{b_{n}}{a} S_{j}^{-2 \beta}\right) \\
& w_{j}^{n}:=\frac{\Delta S^{2}}{12}\left(-\frac{4(\beta+1)}{S_{j}}+\frac{b_{n}}{a} S_{j}^{-2 \beta-1}\right)
\end{align*}
$$

then, insert $g\left(S_{j}, \tau_{n}\right)$ in it

$$
\begin{aligned}
& \gamma_{j}^{n} \delta_{S}^{2} V\left(S_{j}, \tau_{n}\right)+\xi_{j}^{n} \delta_{S} V\left(S_{j}, \tau_{n}\right)=\theta_{j}^{n}\left[{ }_{0}^{C} D_{\tau}^{\alpha} V\left(S_{j}, \tau_{n}\right)+c_{n} V\left(S_{j}, \tau_{n}\right)\right] \\
& \quad+\left[w_{j}^{n} \delta_{S}+\frac{\Delta S^{2}}{12} \delta_{S}^{2}\right]\left[{ }_{0}^{C} D_{\tau}^{\alpha} V\left(S_{j}, \tau_{n}\right)+c_{n} V\left(S_{j}, \tau_{n}\right)\right]+O\left(\Delta S^{4}\right)
\end{aligned}
$$

Now, substitute $\delta_{S}^{2} V\left(S_{j}, \tau_{n}\right), \delta_{S} V\left(S_{j}, \tau_{n}\right)$ and (3.1) in (3.6)

$$
\begin{aligned}
& \gamma_{j}^{n} \frac{V_{j-1}^{n}-2 V_{j}^{n}+V_{j+1}^{n}}{\Delta S^{2}}+\xi_{j}^{n} \frac{V_{j+1}^{n}-V_{j-1}^{n}}{2 \Delta S} \\
& =\theta_{j}^{n}\left[\varphi \sum_{k=1}^{n} \psi_{k}\left(V_{j}^{n-k+1}-V_{j}^{n-k}\right)+c_{n} V\left(S_{j}, \tau_{n}\right)\right] \\
& \quad+\left[w_{j}^{n} \delta_{S}+\frac{\Delta S^{2}}{12} \delta_{S}^{2}\right]\left[\varphi \sum_{k=1}^{n} \psi_{k}\left(V_{j}^{n-k+1}-V_{j}^{n-k}\right)+c_{n} V_{j}^{n}\right]+R_{j}^{n}
\end{aligned}
$$

where the evaluation $R_{j}^{n}$ holds in $\left|R_{j}^{n}\right| \leq C\left(\Delta \tau^{2-\alpha}+\Delta S^{4}\right)$. Again, by substituting $\delta_{S}^{2} V\left(S_{j}, \tau_{n}\right)$ and $\delta_{S} V\left(S_{j}, \tau_{n}\right)$ on the right hand side above formula and rearranging, we obtain

$$
\begin{align*}
& {\left[\frac{\gamma_{j}^{n}}{\Delta S^{2}}-\frac{\xi_{j}^{n}}{2 \Delta S}+\frac{w_{j}^{n}}{2 \Delta S}\left(\varphi+c_{n}\right)-\frac{1}{12}\left(\varphi+c_{n}\right)\right] V_{j-1}^{n}} \\
& \quad-\left[\frac{2 \gamma_{j}^{n}}{\Delta S^{2}}+c_{n} \theta_{j}^{n}+\varphi \theta_{j}^{n}-\frac{1}{6}\left(\varphi+c_{n}\right)\right] V_{j}^{n} \\
& \quad+\left[\frac{\gamma_{j}^{n}}{\Delta S^{2}}+\frac{\xi_{j}^{n}}{2 \Delta S}-\frac{w_{j}^{n}}{2 \Delta S}\left(\varphi+c_{n}\right)-\frac{1}{12}\left(\varphi+c_{n}\right)\right] V_{j+1}^{n} \\
& =-\theta_{j}^{n} \varphi V_{j}^{n-1}+\theta_{j}^{n} \varphi \sum_{k=2}^{n} \psi_{k}\left(V_{j}^{n-k+1}-V_{j}^{n-k}\right) \\
& \quad-\frac{w_{j}^{n}}{2 \Delta S} \varphi\left(V_{j+1}^{n-1}-V_{j-1}^{n-1}\right) \\
& \quad+\frac{w_{j}^{n}}{2 \Delta S} \varphi \sum_{k=2}^{n} \psi_{k}\left(V_{j+1}^{n-k+1}-V_{j-1}^{n-k+1}-V_{j+1}^{n-k}+V_{j-1}^{n-k}\right)  \tag{3.7}\\
& \quad-\frac{1}{12} \varphi\left(V_{j-1}^{n-1}-2 V_{j}^{n-1}+V_{j+1}^{n-1}\right) \\
& \quad+\frac{1}{12} \varphi \sum_{k=2}^{n} \psi_{k}\left(V_{j-1}^{n-k+1}-2 V_{j}^{n-k+1}+V_{j+1}^{n-k+1}-V_{j-1}^{n-k}\right. \\
& \left.\quad+2 V_{j}^{n-k}-V_{j+1}^{n-k}\right)+R_{j}^{n} .
\end{align*}
$$

Finally, with eliminating $R_{j}^{n}$ and rearranging, we get the following compact difference scheme

$$
\begin{aligned}
& {\left[\frac{\gamma_{j}^{n}}{\Delta S^{2}}-\frac{\xi_{j}^{n}}{2 \Delta S}+\left(\varphi+c_{n}\right)\left(\frac{w_{j}^{n}}{2 \Delta S}-\frac{1}{12}\right)\right] \widetilde{V}_{j-1}^{n}} \\
& \quad-\left[\frac{2 \gamma_{j}^{n}}{\Delta S^{2}}+c_{n} \theta_{j}^{n}+\left(\varphi+c_{n}\right)\left(\theta_{j}^{n}-\frac{1}{6}\right)\right] \widetilde{V}_{j}^{n} \\
& \quad+\left[\frac{\gamma_{j}^{n}}{\Delta S^{2}}+\frac{\xi_{j}^{n}}{2 \Delta S}-\left(\varphi+c_{n}\right)\left(\frac{w_{j}^{n}}{2 \Delta S}+\frac{1}{12}\right)\right] \widetilde{V}_{j+1}^{n}
\end{aligned}
$$

$$
\begin{align*}
= & \theta_{j}^{n} \varphi\left[\sum_{k=1}^{n-1}\left(\psi_{k+1}-\psi_{k}\right) \widetilde{V}_{j}^{n-k}-\psi_{n} \widetilde{V}_{j}^{0}\right] \\
& +\frac{w_{j}^{n}}{2 \Delta S} \varphi\left[\sum_{k=1}^{n-1}\left(\psi_{k+1}-\psi_{k}\right)\left(\widetilde{V}_{j+1}^{n-k}-\widetilde{V}_{j-1}^{n-k}\right)-\psi_{n}\left(\widetilde{V}_{j+1}^{0}-\widetilde{V}_{j-1}^{0}\right)\right] \\
& +\frac{\varphi}{12}\left[\sum_{k=1}^{n-1}\left(\psi_{k+1}-\psi_{k}\right)\left(\widetilde{V}_{j-1}^{n-k}-2 \widetilde{V}_{j}^{n-k}+\widetilde{V}_{j+1}^{n-k}\right)\right.  \tag{3.8}\\
& \left.-\psi_{n}\left(\widetilde{V}_{j-1}^{0}-2 \widetilde{V}_{j}^{0}+\widetilde{V}_{j+1}^{0}\right)\right],
\end{align*}
$$

where $\widetilde{V}$ is the exact solution of the compact difference scheme and $V$ is the exact solution of the differential equation.

An index $J\left(\tau_{n}\right)$ (for $\left.n=0,1, \ldots, M\right)$ is found such that (similar to [53])

$$
\begin{aligned}
& \frac{1}{2} \delta^{2} S_{j}^{2 \beta+2} \frac{\partial^{2} V\left(S_{j}, \tau_{n}\right)}{\partial S^{2}}+\left(r\left(\tau_{n}\right)-D\left(\tau_{n}\right)\right) S_{j} \frac{\partial V\left(S_{j}, \tau_{n}\right)}{\partial S} \\
& \quad-r\left(\tau_{n}\right) V\left(S_{j}, \tau_{n}\right)-{ }_{0}^{C} D_{\tau}^{\alpha} V\left(S_{j}, \tau_{n}\right)=0,
\end{aligned}
$$

where $V\left(S_{j}, \tau_{n}\right)>\max \left(S_{j}-K, 0\right)$ for $j=0,1, \ldots, J\left(\tau_{n}\right)$, and

$$
\begin{aligned}
& \frac{1}{2} \delta^{2} S_{j}^{2 \beta+2} \frac{\partial^{2} V\left(S_{j}, \tau_{n}\right)}{\partial S^{2}}+\left(r\left(\tau_{n}\right)-D\left(\tau_{n}\right)\right) S_{j} \frac{\partial V\left(S_{j}, \tau_{n}\right)}{\partial S} \\
& \quad-r\left(\tau_{n}\right) V\left(S_{j}, \tau_{n}\right)-{ }_{0}^{C} D_{\tau}^{\alpha} V\left(S_{j}, \tau_{n}\right)<0,
\end{aligned}
$$

where $V\left(S_{j}, \tau_{n}\right)=\max \left(S_{j}-K, 0\right)$ for $j=J\left(\tau_{n}\right)+1, J\left(\tau_{n}\right)+2, \ldots, N$. Therefore, American call option prices are obtained from (3.8) for $j=0,1, \ldots, J\left(\tau_{n}\right)$ and are equal $S_{j}-K$ for $j=J\left(\tau_{n}\right)+1, \ldots, N$.

## 4. Stability and convergence analysis

In this section, we investigate that the compact difference scheme is solvable, unconditionally stable and convergent.

### 4.1. Solvability.

Theorem 4.1. The compact difference scheme (3.8) has a unique solution.
Proof. Matrix form of compact difference scheme (3.8) can be briefly written as $A^{n} \widetilde{V}^{n}=d_{n-1}$, where $d_{n-1}$ depends only $\widetilde{V}^{n-1}, \widetilde{V}^{n-2}, \ldots \widetilde{V}^{0}$. The tridiagonal coefficient matrix $A^{n}=\left(a_{i j}^{n}\right)$ from compact difference scheme (3.8) is strictly diagonally dominant since $\left|a_{i i}^{n}\right|>\sum_{j \neq i}\left|a_{i j}^{n}\right|$, where

$$
\begin{aligned}
& \left|a_{i i}^{n}\right|=\frac{2 \gamma_{i}^{n}}{\Delta S^{2}}+c_{n} \theta_{i}^{n}+\left(\varphi+c_{n}\right)\left(\theta_{i}^{n}-\frac{1}{6}\right) \\
& \sum_{j \neq i}\left|a_{i j}^{n}\right|=\frac{2 \gamma_{j}^{n}}{\Delta S^{2}}-\frac{1}{6}\left(\varphi+c_{n}\right), \quad n=1,2, \ldots, M
\end{aligned}
$$

For each $n$, the coefficient matrix $A^{n}$ is nonsingular, so the compact difference scheme (3.8) is uniquely solvable.
4.2. Stability. In this part, we prove that the compact difference scheme (3.8) is unconditionally stable using the Fourier analysis [45, 12, 24]. Suppose $\widehat{V}_{j}^{n}$ be a numerical solution of compact difference scheme (3.8). Let

$$
\varepsilon_{j}^{n}=\widetilde{V}_{j}^{n}-\widehat{V}_{j}^{n}, \quad j=0,1, \ldots, N ; \quad n=0,1, \ldots, M
$$

then, $\varepsilon_{j}^{n}$ satisfies in the following equations

$$
\begin{align*}
& {\left[\frac{\gamma_{j}^{n}}{\Delta S^{2}}-\frac{\xi_{j}^{n}}{2 \Delta S}+\left(\varphi+c_{n}\right)\left(\frac{w_{j}^{n}}{2 \Delta S}-\frac{1}{12}\right)\right] \varepsilon_{j-1}^{n}} \\
& \quad-\left[\frac{2 \gamma_{j}^{n}}{\Delta S^{2}}+c_{n} \theta_{j}^{n}+\left(\varphi+c_{n}\right)\left(\theta_{j}^{n}-\frac{1}{6}\right)\right] \varepsilon_{j}^{n} \\
& \quad+\left[\frac{\gamma_{j}^{n}}{\Delta S^{2}}+\frac{\xi_{j}^{n}}{2 \Delta S}-\left(\varphi+c_{n}\right)\left(\frac{w_{j}^{n}}{2 \Delta S}+\frac{1}{12}\right)\right] \varepsilon_{j+1}^{n} \\
& =\theta_{j}^{n} \varphi\left[\sum_{k=1}^{n-1}\left(\psi_{k+1}-\psi_{k}\right) \varepsilon_{j}^{n-k}-\psi_{n} \varepsilon_{j}^{0}\right] \\
& \quad+\frac{w_{j}^{n}}{2 \Delta S} \varphi\left[\sum_{k=1}^{n-1}\left(\psi_{k+1}-\psi_{k}\right)\left(\varepsilon_{j+1}^{n-k}-\varepsilon_{j-1}^{n-k}\right)-\psi_{n}\left(\varepsilon_{j+1}^{0}-\varepsilon_{j-1}^{0}\right)\right]  \tag{4.1}\\
& \quad+\frac{\varphi}{12}\left[\sum_{k=1}^{n-1}\left(\psi_{k+1}-\psi_{k}\right)\left(\varepsilon_{j-1}^{n-k}-2 \varepsilon_{j}^{n-k}+\varepsilon_{j+1}^{n-k}\right)\right. \\
& \left.\quad-\psi_{n}\left(\varepsilon_{j-1}^{0}-2 \varepsilon_{j}^{0}+\varepsilon_{j+1}^{0}\right)\right],
\end{align*}
$$

and $\varepsilon_{0}^{n}=\varepsilon_{J\left(\tau_{n}\right)+1}^{n}=\cdots=\varepsilon_{N}^{n}=0$.
Now, we define the following grid function

$$
\varepsilon^{n}(S)=\left\{\begin{array}{cc}
\varepsilon_{j}^{n}, & S \in\left(S_{j}-\frac{\Delta S}{2}, S_{j}+\frac{\Delta S}{2}\right] \\
0, & S \in\left[0, \frac{\Delta S}{2}\right] \bigcup\left(S_{J\left(\tau_{n}\right)+1}-\frac{\Delta S}{2}, S_{J\left(\tau_{n}\right)+1}+\frac{\Delta S}{2}\right] \\
\bigcup \cdots \bigcup\left(S_{\max }-\frac{\Delta S}{2}, S_{\max }\right]
\end{array}\right.
$$

and make a Fourier series extension for it with the period $L=S_{\max }$ as follows

$$
\begin{aligned}
& \varepsilon^{n}(S)=\sum_{j=-\infty}^{+\infty} \varsigma_{j}^{n} e^{i \frac{2 \pi j S}{L}} \quad\left(i^{2}=-1\right), \quad n=0,1, \ldots, M, \\
& \varsigma_{j}^{n}=\frac{1}{L} \int_{0}^{L} \varepsilon^{n}(S) e^{i \frac{2 \pi j S}{L}} d S, \quad j \in \mathbb{Z} .
\end{aligned}
$$

We let $\varepsilon^{n}=\left(\varepsilon_{1}^{n}, \varepsilon_{2}^{n}, \ldots, \varepsilon_{J\left(\tau_{n}\right)}^{n}\right)^{t}$, and define the following norm

$$
\left\|\varepsilon^{n}\right\|_{2}^{2}=\sum_{j=1}^{J\left(\tau_{n}\right)} \Delta S\left|\varepsilon_{j}^{n}\right|^{2}=\int_{0}^{L}\left|\varepsilon^{n}(S)\right|^{2} d S=\left\|\varepsilon^{n}(S)\right\|_{L^{2}}^{2}, \quad n=0,1, \ldots, M
$$

In addition, by using the Parseval equality

$$
\left\|\varepsilon^{n}(S)\right\|_{L^{2}}^{2}=L \sum_{j=-\infty}^{+\infty}\left|\varsigma_{j}^{n}\right|^{2}, \quad n=0,1, \ldots, M
$$

we obtain

$$
\begin{equation*}
\left\|\varepsilon^{n}\right\|_{2}^{2}=\sum_{j=1}^{J\left(\tau_{n}\right)} \Delta S\left|\varepsilon_{j}^{n}\right|^{2}=L \sum_{j=-\infty}^{+\infty}\left|\varsigma_{j}^{n}\right|^{2}, \quad n=0,1, \ldots, M \tag{4.2}
\end{equation*}
$$

According to the above analysis and $S_{j}=j \Delta S$, we assume that the solution of (3.8) has the form as follows

$$
\varepsilon_{j}^{n}=\varsigma^{n} e^{i q j \Delta S}, \quad q=\frac{2 \pi l}{L}, \quad l \in \mathbb{Z}
$$

With substituting the above formula into (4.1), we obtain

$$
\begin{aligned}
& {\left[\frac{\gamma_{j}^{n}}{\Delta S^{2}}-\frac{\xi_{j}^{n}}{2 \Delta S}+\left(\varphi+c_{n}\right)\left(\frac{w_{j}^{n}}{2 \Delta S}-\frac{1}{12}\right)\right] \varsigma^{n} e^{-i q \Delta S}} \\
& \quad-\left[\frac{2 \gamma_{j}^{n}}{\Delta S^{2}}+c_{n} \theta_{j}^{n}+\left(\varphi+c_{n}\right)\left(\theta_{j}^{n}-\frac{1}{6}\right)\right] \varsigma^{n} \\
& \quad+\left[\frac{\gamma_{j}^{n}}{\Delta S^{2}}+\frac{\xi_{j}^{n}}{2 \Delta S}-\left(\varphi+c_{n}\right)\left(\frac{w_{j}^{n}}{2 \Delta S}+\frac{1}{12}\right)\right] \varsigma^{n} e^{i q \Delta S} \\
& =\theta_{j}^{n} \varphi\left[\sum_{k=1}^{n-1}\left(\psi_{k+1}-\psi_{k}\right) \varsigma^{n-k}-\psi_{n} \varsigma^{0}\right] \\
& \quad+\frac{w_{j}^{n}}{2 \Delta S} \varphi\left[\sum_{k=1}^{n-1}\left(\psi_{k+1}-\psi_{k}\right) \varsigma^{n-k}\left(e^{i q \Delta S}-e^{-i q \Delta S}\right)\right. \\
& \left.\quad-\psi_{n} \varsigma^{0}\left(e^{i q \Delta S}-e^{-i q \Delta S}\right)\right] \\
& \quad+\frac{\varphi}{12}\left[\sum_{k=1}^{n-1}\left(\psi_{k+1}-\psi_{k}\right) \varsigma^{n-k}\left(e^{-i q \Delta S}-2+e^{i q \Delta S}\right)\right. \\
& \left.\quad-\psi_{n} \varsigma^{0}\left(e^{-i q \Delta S}-2+e^{i q \Delta S}\right)\right] .
\end{aligned}
$$

By using $\sin ^{2}\left(\frac{q \Delta S}{2}\right)=-\frac{1}{4}\left(e^{i q \Delta S}-2+e^{-i q \Delta S}\right)$, the above relation reduces to

$$
\begin{aligned}
& {\left[\left(-\frac{4 \gamma_{j}^{n}}{\Delta S^{2}}+\frac{1}{3}\left(\varphi+c_{n}\right)\right) \sin ^{2}\left(\frac{q \Delta S}{2}\right)-\left(\varphi+2 c_{n}\right) \theta_{j}^{n}\right.} \\
& \left.\quad+i\left(\frac{\xi_{j}^{n}}{\Delta S}-\left(\varphi+c_{n}\right) \frac{w_{j}^{n}}{\Delta S}\right) \sin (q \Delta S)\right] \varsigma^{n} \\
& =\left[\theta_{j}^{n} \varphi+i \frac{w_{j}^{n}}{\Delta S} \varphi \sin (q \Delta S)-\frac{\varphi}{3} \sin ^{2}\left(\frac{q \Delta S}{2}\right)\right] \\
& \quad \times\left[\sum_{k=1}^{n-1}\left(\psi_{k+1}-\psi_{k}\right) \varsigma^{n-k}-\psi_{n} \varsigma^{0}\right] .
\end{aligned}
$$

This yields

$$
\begin{align*}
\varsigma^{n}= & {\left[\theta_{j}^{n} \varphi+i \frac{w_{j}^{n}}{\Delta S} \varphi \sin (q \Delta S)-\frac{\varphi}{3} \sin ^{2}\left(\frac{q \Delta S}{2}\right)\right] / } \\
& {\left[\left(-\frac{4 \gamma_{j}^{n}}{\Delta S^{2}}+\frac{\varphi+c_{n}}{3}\right) \sin ^{2}\left(\frac{q \Delta S}{2}\right)-\left(\varphi+2 c_{n}\right) \theta_{j}^{n}\right.} \\
& \left.+i\left(\frac{\xi_{j}^{n}}{\Delta S}-\left(\varphi+c_{n}\right) \frac{w_{j}^{n}}{\Delta S}\right) \sin (q \Delta S)\right]  \tag{4.3}\\
& \times\left[\sum_{k=1}^{n-1}\left(\psi_{k+1}-\psi_{k}\right) \varsigma^{n-k}-\psi_{n} \varsigma^{0}\right] .
\end{align*}
$$

Lemma 4.2. The coefficients $\psi_{n}$ satisfy
(I) $\psi_{n}>0, \quad n=1,2, \ldots$
(II) $1=\psi_{1}>\psi_{2}>\ldots>\psi_{n+1}, \quad \psi_{n+1} \rightarrow 0, \quad$ as $n \rightarrow+\infty$.
(III) $\sum_{k=1}^{n}\left(\psi_{k}-\psi_{k+1}\right)+\psi_{n+1}=1$.

Proof. (I)-(III) are clearly established.
Lemma 4.3. The following inequality is established

$$
\begin{aligned}
& \left\lvert\,\left[\theta_{j}^{n} \varphi+i \frac{w_{j}^{n}}{\Delta S} \varphi \sin (q \Delta S)-\frac{\varphi}{3} \sin ^{2}\left(\frac{q \Delta S}{2}\right)\right] /\right. \\
& \quad\left[\left(-\frac{4 \gamma_{j}^{n}}{\Delta S^{2}}+\frac{\varphi+c_{n}}{3}\right) \sin ^{2}\left(\frac{q \Delta S}{2}\right)-\left(\varphi+2 c_{n}\right) \theta_{j}^{n}\right. \\
& \left.\quad+i\left(\frac{\xi_{j}^{n}}{\Delta S}-\left(\varphi+c_{n}\right) \frac{w_{j}^{n}}{\Delta S}\right) \sin (q \Delta S)\right] \mid \leq 1
\end{aligned}
$$

Proof. See [14].
Lemma 4.4. Suppose that $\varsigma^{n}(n=1,2, \cdots, N)$ is the solutions of (4.3), we have $\left|\varsigma^{n}\right| \leq\left|\varsigma^{0}\right|$.

Proof. We prove it by using the mathematical induction. For $n=1,(4.3)$ becomes

$$
\begin{aligned}
\varsigma^{1}= & {\left[\left(\theta_{j}^{n} \varphi+i \frac{w_{j}^{n}}{\Delta S} \varphi \sin (q \Delta S)-\frac{\varphi}{3} \sin ^{2}\left(\frac{q \Delta S}{2}\right)\right) \varsigma^{0}\right] / } \\
& {\left[\left(-\frac{4 \gamma_{j}^{n}}{\Delta S^{2}}+\frac{\varphi+c_{n}}{3}\right) \sin ^{2}\left(\frac{q \Delta S}{2}\right)-\left(\varphi+2 c_{n}\right) \theta_{j}^{n}\right.} \\
& \left.+i\left(\frac{\xi_{j}^{n}}{\Delta S}-\left(\varphi+c_{n}\right) \frac{w_{j}^{n}}{\Delta S}\right) \sin (q \Delta S)\right] .
\end{aligned}
$$

Noticing that $\psi_{1}=1$, according to Lemma 4.3, $\left|\varsigma^{1}\right| \leq\left|\varsigma^{0}\right|$. Suppose that $\left|\varsigma^{k}\right| \leq\left|\varsigma^{0}\right|$, $k=2,3, \ldots, n-1$, and prove $\left|\varsigma^{n}\right| \leq\left|\varsigma^{0}\right|$. By applying Lemma 4.2, Lemma 4.3 and the relation (4.3) for $n \geq 2$, we get

$$
\begin{aligned}
&\left|\varsigma^{n}\right|= \mid \\
& {\left[\left(\theta_{j}^{n} \varphi+i \frac{w_{j}^{n}}{\Delta S} \varphi \sin (q \Delta S)-\frac{\varphi}{3} \sin ^{2}\left(\frac{q \Delta S}{2}\right)\right] /\right.} \\
& \quad\left[\left(-\frac{4 \gamma_{j}^{n}}{\Delta S^{2}}+\frac{\varphi+c_{n}}{3}\right) \sin ^{2}\left(\frac{q \Delta S}{2}\right)-\left(\varphi+2 c_{n}\right) \theta_{j}^{n}\right. \\
&\left.\left.\times \left\lvert\, \sum_{k=1}^{\Delta S}-\left(\varphi+c_{n}\right) \frac{w_{j}^{n}}{\Delta S}\right.\right) \sin (q \Delta S)\right] \mid \\
&\left.\leq \mid \psi_{k+1}^{n-1}-\psi_{k}\right) \varsigma^{n-k}-\psi_{n} \varsigma^{0} \mid \\
& \leq\left(\psi_{k+1}^{n-1}-\psi_{k}\right) \varsigma^{n-k}-\psi_{n} \varsigma^{0} \mid \\
& \leq \sum_{k=1}^{n-1}\left(\psi_{k+1}-\psi_{k}\right) \varsigma^{n-k}\left|+\psi_{n}\right| \varsigma^{0} \mid \\
& \leq \sum_{k=1}^{n-1}\left(\psi_{k}-\psi_{k+1}\right)\left|\varsigma^{n-k}\right|+\psi_{n}\left|\varsigma^{0}\right| \\
& \leq\left(\sum_{k=1}^{n-1}\left(\psi_{k}-\psi_{k+1}\right)+\psi_{n}\right)\left|\varsigma^{0}\right|=\left|\varsigma^{0}\right|
\end{aligned}
$$

thus $\left|\varsigma^{n}\right| \leq\left|\varsigma^{0}\right|$. The proof is finished.
Theorem 4.5. The compact difference scheme (3.8) is unconditionally stable.
Proof. According to Lemma 4.4 and formula (4.2), we derive

$$
\left\|\varepsilon^{n}\right\|_{2}^{2}=L \sum_{j=-\infty}^{+\infty}\left|\varsigma_{j}^{n}\right|^{2} \leq L \sum_{j=-\infty}^{+\infty}\left|\varsigma_{j}^{0}\right|^{2}=\left\|\varepsilon^{0}\right\|_{2}^{2}
$$

This yields $\left\|\varepsilon^{n}\right\|_{2} \leq\left\|\varepsilon^{0}\right\|_{2}$ for $n=1,2, \ldots, M$. Thus, the compact difference scheme (3.8) is unconditionally stable.
4.3. Convergence. Now, we investigate the convergence of the compact difference scheme (3.8). We suppose that $V_{j}^{n}$ is the exact solution of (2.5) and $\widetilde{V}_{j}^{n}$ is the exact solution of (3.8). Let

$$
\left\{\begin{array}{l}
E_{j}^{n}=V_{j}^{n}-\widetilde{V}_{j}^{n}, \quad j=0,1, \ldots, N ; \quad n=0,1, \ldots, M  \tag{4.4}\\
R_{j}^{n}=O\left(\Delta \tau^{2-\alpha}+\Delta S^{4}\right)
\end{array}\right.
$$

Subtracting (3.7) from (3.8) achieve

$$
\begin{align*}
& {\left[\frac{\gamma_{j}^{n}}{\Delta S^{2}}-\frac{\xi_{j}^{n}}{2 \Delta S}+\left(\varphi+c_{n}\right)\left(\frac{w_{j}^{n}}{2 \Delta S}-\frac{1}{12}\right)\right] E_{j-1}^{n}} \\
& \quad-\left[\frac{2 \gamma_{j}^{n}}{\Delta S^{2}}+c_{n} \theta_{j}^{n}+\left(\varphi+c_{n}\right)\left(\theta_{j}^{n}-\frac{1}{6}\right)\right] E_{j}^{n} \\
& \quad+\left[\frac{\gamma_{j}^{n}}{\Delta S^{2}}+\frac{\xi_{j}^{n}}{2 \Delta S}-\left(\varphi+c_{n}\right)\left(\frac{w_{j}^{n}}{2 \Delta S}+\frac{1}{12}\right)\right] E_{j+1}^{n} \\
& =\theta_{j}^{n} \varphi\left[\sum_{k=1}^{n-1}\left(\psi_{k+1}-\psi_{k}\right) E_{j}^{n-k}-\psi_{n} E_{j}^{0}\right] \\
& \quad+\frac{w_{j}^{n}}{2 \Delta S} \varphi\left[\sum_{k=1}^{n-1}\left(\psi_{k+1}-\psi_{k}\right)\left(E_{j+1}^{n-k}-E_{j-1}^{n-k}\right)-\psi_{n}\left(E_{j+1}^{0}-E_{j-1}^{0}\right)\right]  \tag{4.5}\\
& \quad+\frac{\varphi}{12}\left[\sum_{k=1}^{n-1}\left(\psi_{k+1}-\psi_{k}\right)\left(E_{j-1}^{n-k}-2 E_{j}^{n-k}+E_{j+1}^{n-k}\right)\right. \\
& \left.\quad-\psi_{n}\left(E_{j-1}^{0}-2 E_{j}^{0}+E_{j+1}^{0}\right)\right]+R_{j}^{n}
\end{align*}
$$

with the initial-boundary conditions

$$
\left\{\begin{array}{l}
E_{j}^{0}=0, \quad j=0,1, \ldots, N \\
E_{0}^{n}=E_{J\left(\tau_{n}\right)+1}^{n}=\cdots=E_{N}^{n}=0, \quad n=0,1, \ldots, M
\end{array}\right.
$$

Similar to the stability proof as above, we define the following grid functions

$$
E^{n}(S)=\left\{\begin{array}{c}
E_{j}^{n}, \quad S \in\left(S_{j}-\frac{\Delta S}{2}, S_{j}+\frac{\Delta S}{2}\right] \\
0, \quad S \in\left[0, \frac{\Delta S}{2}\right] \cup\left(S_{J\left(\tau_{n}\right)+1}-\frac{\Delta S}{2}, S_{J\left(\tau_{n}\right)+1}+\frac{\Delta S}{2}\right] \\
\bigcup \cdots \bigcup\left(S_{\max }-\frac{\Delta S}{2}, S_{\max }\right],
\end{array}\right.
$$

and

$$
R^{n}(S)=\left\{\begin{array}{c}
R_{j}^{n}, \quad S \in\left(S_{j}-\frac{\Delta S}{2}, S_{j}+\frac{\Delta S}{2}\right] \\
0, \quad S \in\left[0, \frac{\Delta S}{2}\right] \cup\left(S_{J\left(\tau_{n}\right)+1}-\frac{\Delta S}{2}, S_{J\left(\tau_{n}\right)+1}+\frac{\Delta S}{2}\right] \\
\bigcup \cdots \bigcup\left(S_{\max }-\frac{\Delta S}{2}, S_{\max }\right]
\end{array}\right.
$$

Therefore, $E^{n}(S)$ and $R^{n}(S)$ have the Fourier series expansions as follow, respectively,

$$
\begin{aligned}
& E^{n}(S)=\sum_{j=-\infty}^{+\infty} \vartheta_{j}^{n} e^{i \frac{2 \pi j S}{L}} \\
& R^{n}(S)=\sum_{j=-\infty}^{+\infty} \nu_{j}^{n} e^{i \frac{2 \pi j S}{L}}, \quad\left(i^{2}=-1\right), \quad n=0,1, \ldots, M
\end{aligned}
$$

where $L=S_{\text {max }}$ and

$$
\vartheta_{j}^{n}=\frac{1}{L} \int_{0}^{L} E^{n}(S) e^{i \frac{2 \pi j S}{L}} d S, \quad \nu_{j}^{n}=\frac{1}{L} \int_{0}^{L} R^{n}(S) e^{i \frac{2 \pi j S}{L}} d S, \quad j \in \mathbb{Z}
$$

Let $E^{n}=\left(E_{1}^{n}, E_{2}^{n}, \ldots, E_{J\left(\tau_{n}\right)}^{n}\right)^{t}, R^{n}=\left(R_{1}^{n}, R_{2}^{n}, \ldots, R_{J\left(\tau_{n}\right)}^{n}\right)^{t}$ and define their corresponding norms as follows, respectively

$$
\begin{aligned}
\left\|E^{n}\right\|_{2}^{2}=\sum_{j=1}^{J\left(\tau_{n}\right)} \Delta S\left|E_{j}^{n}\right|^{2}=\int_{0}^{L}\left|E^{n}(S)\right|^{2} d S= & \left\|E^{n}(S)\right\|_{L^{2}}^{2} \\
& n=0,1, \ldots, M
\end{aligned}
$$

and

$$
\begin{align*}
\left\|R^{n}\right\|_{2}^{2}=\sum_{j=1}^{J\left(\tau_{n}\right)} \Delta S\left|R_{j}^{n}\right|^{2}=\int_{0}^{L}\left|R^{n}(S)\right|^{2} d S= & \left\|R^{n}(S)\right\|_{L^{2}}^{2}  \tag{4.6}\\
& n=0,1, \ldots, M
\end{align*}
$$

Applying Parseval equality leads to

$$
\begin{align*}
& \left\|E^{n}\right\|_{2}^{2}=\sum_{j=1}^{J\left(\tau_{n}\right)} \Delta S\left|E_{j}^{n}\right|^{2}=L \sum_{j=-\infty}^{+\infty}\left|\vartheta_{j}^{n}\right|^{2}, \quad n=0,1, \ldots, M  \tag{4.7a}\\
& \left\|R^{n}\right\|_{2}^{2}=\sum_{j=1}^{J\left(\tau_{n}\right)} \Delta S\left|R_{j}^{n}\right|^{2}=L \sum_{j=-\infty}^{+\infty}\left|\nu_{j}^{n}\right|^{2}, \quad n=0,1, \ldots, M \tag{4.7b}
\end{align*}
$$

According to the stability analysis and $S_{j}=j \Delta S$, we suppose that the solution of (4.5) has the form as follows

$$
E_{j}^{n}=\vartheta^{n} e^{i q j \Delta S}, \quad R_{j}^{n}=\nu^{n} e^{i q j \Delta S}, \quad q=\frac{2 \pi l}{L}, \quad l \in \mathbb{Z}
$$

Replacing the above relations into (4.5), we have

$$
\begin{aligned}
& {\left[\frac{\gamma_{j}^{n}}{\Delta S^{2}}-\frac{\xi_{j}^{n}}{2 \Delta S}+\left(\varphi+c_{n}\right)\left(\frac{w_{j}^{n}}{2 \Delta S}-\frac{1}{12}\right)\right] \vartheta^{n} e^{-i q \Delta S}} \\
& \quad-\left[\frac{2 \gamma_{j}^{n}}{\Delta S^{2}}+c_{n} \theta_{j}^{n}+\left(\varphi+c_{n}\right)\left(\theta_{j}^{n}-\frac{1}{6}\right)\right] \vartheta^{n} \\
& \quad+\left[\frac{\gamma_{j}^{n}}{\Delta S^{2}}+\frac{\xi_{j}^{n}}{2 \Delta S}-\left(\varphi+c_{n}\right)\left(\frac{w_{j}^{n}}{2 \Delta S}+\frac{1}{12}\right)\right] \vartheta^{n} e^{i q \Delta S}
\end{aligned}
$$

$$
\begin{aligned}
= & \theta_{j}^{n} \varphi\left[\sum_{k=1}^{n-1}\left(\psi_{k+1}-\psi_{k}\right) \vartheta^{n-k}-\psi_{n} \vartheta^{0}\right] \\
& +\frac{w_{j}^{n}}{2 \Delta S} \varphi\left[\sum_{k=1}^{n-1}\left(\psi_{k+1}-\psi_{k}\right)\left(e^{i q \Delta S}-e^{-i q \Delta S}\right) \vartheta^{n-k}\right. \\
& \left.-\psi_{n}\left(e^{i q \Delta S}-e^{-i q \Delta S}\right) \vartheta^{0}\right] \\
& +\frac{\varphi}{12}\left[\sum_{k=1}^{n-1}\left(\psi_{k+1}-\psi_{k}\right)\left(e^{-i q \Delta S}-2+e^{i q \Delta S}\right) \vartheta^{n-k}\right. \\
& \left.-\psi_{n}\left(e^{-i q \Delta S}-2+e^{i q \Delta S}\right) \vartheta^{0}\right]+\nu^{n} .
\end{aligned}
$$

Applying $\sin ^{2}\left(\frac{q \Delta S}{2}\right)=-\frac{1}{4}\left(e^{i q \Delta S}-2+e^{-i q \Delta S}\right)$, the above relation turns to

$$
\begin{aligned}
& {\left[\left(-\frac{4 \gamma_{j}^{n}}{\Delta S^{2}}+\frac{1}{3}\left(\varphi+c_{n}\right)\right) \sin ^{2}\left(\frac{q \Delta S}{2}\right)-\left(\varphi+2 c_{n}\right) \theta_{j}^{n}\right.} \\
& \left.\quad+i\left(\frac{\xi_{j}^{n}}{\Delta S}-\left(\varphi+c_{n}\right) \frac{w_{j}^{n}}{\Delta S}\right) \sin (q \Delta S)\right] \vartheta^{n} \\
& =\left[\theta_{j}^{n} \varphi+i \frac{w_{j}^{n}}{\Delta S} \varphi \sin (q \Delta S)-\frac{\varphi}{3} \sin ^{2}\left(\frac{q \Delta S}{2}\right)\right] \\
& \quad \times\left[\sum_{k=1}^{n-1}\left(\psi_{k+1}-\psi_{k}\right) \vartheta^{n-k}-\psi_{n} \vartheta^{0}\right]+\nu^{n}
\end{aligned}
$$

in result

$$
\begin{align*}
\vartheta^{n}= & {\left[\theta_{j}^{n} \varphi+i \frac{w_{j}^{n}}{\Delta S} \varphi \sin (q \Delta S)-\frac{\varphi}{3} \sin ^{2}\left(\frac{q \Delta S}{2}\right)\right] / } \\
& {\left[\left(-\frac{4 \gamma_{j}^{n}}{\Delta S^{2}}+\frac{\varphi+c_{n}}{3}\right) \sin ^{2}\left(\frac{q \Delta S}{2}\right)-\left(\varphi+2 c_{n}\right) \theta_{j}^{n}\right.} \\
& \left.+i\left(\frac{\xi_{j}^{n}}{\Delta S}-\left(\varphi+c_{n}\right) \frac{w_{j}^{n}}{\Delta S}\right) \sin (q \Delta S)\right] \times \sum_{k=1}^{n-1}\left(\psi_{k+1}-\psi_{k}\right) \vartheta^{n-k}  \tag{4.8}\\
& +\nu^{n} /\left[\left(-\frac{4 \gamma_{j}^{n}}{\Delta S^{2}}+\frac{\varphi+c_{n}}{3}\right) \sin ^{2}\left(\frac{q \Delta S}{2}\right)-\left(\varphi+2 c_{n}\right) \theta_{j}^{n}\right. \\
& \left.+i\left(\frac{\xi_{j}^{n}}{\Delta S}-\left(\varphi+c_{n}\right) \frac{w_{j}^{n}}{\Delta S}\right) \sin (q \Delta S)\right]
\end{align*}
$$

Noticing that $\vartheta^{0}=0$. From (4.4) and (4.6), we obtain ( $C_{1}$ is constant)

$$
\begin{array}{r}
\left\|R^{n}\right\|_{2} \leq \sqrt{N \Delta S} C_{1}\left(\Delta \tau^{2-\alpha}+\Delta S^{4}\right)=\sqrt{L} C_{1}\left(\Delta \tau^{2-\alpha}+\Delta S^{4}\right) \\
n=0,1, \ldots, M \tag{4.9}
\end{array}
$$

Because of the convergence of series in the right hand side of (4.7b), there is a positive constant $C_{2}$ such that

$$
\begin{equation*}
\left|\nu^{n}\right| \equiv\left|\nu_{j}^{n}\right| \leq C_{2} \Delta \tau\left|\nu_{j}^{1}\right| \equiv C_{2} \Delta \tau\left|\nu^{1}\right|, \quad n=1,2, \ldots, M \tag{4.10}
\end{equation*}
$$

Lemma 4.6. The following relationship is established

$$
\begin{aligned}
& 1 /\left[\left(\left[-\frac{4 \gamma_{j}^{n}}{\Delta S^{2}}+\frac{\varphi+c_{n}}{3}\right] \sin ^{2}\left(\frac{q \Delta S}{2}\right)-\left(\varphi+2 c_{n}\right) \theta_{j}^{n}\right)^{2}\right. \\
& \left.+\left(\frac{\xi_{j}^{n}}{\Delta S}-\left(\varphi+c_{n}\right) \frac{w_{j}^{n}}{\Delta S}\right)^{2} \sin ^{2}(q \Delta S)\right] \leq 9
\end{aligned}
$$

Lemma 4.7. Suppose that $\vartheta^{n}$ is a solution of (4.8), then there is a positive constant $C_{3}$ such that

$$
\left|\vartheta^{n}\right| \leq C_{3}(1+3 \Delta \tau)^{n}\left|\nu^{1}\right|, \quad n=1,2, \ldots, M
$$

Proof. We show the proof by using the mathematical induction. From (4.8), (4.10) and (4.6), we obtain

$$
\begin{aligned}
\left|\vartheta^{1}\right|^{2} \leq & \left|\nu^{1}\right|^{2} /\left[\left(\left[-\frac{4 \gamma_{j}^{n}}{\Delta S^{2}}+\frac{\varphi+c_{n}}{3}\right] \sin ^{2}\left(\frac{q \Delta S}{2}\right)-\left(\varphi+2 c_{n}\right) \theta_{j}^{n}\right)^{2}\right. \\
& \left.+\left(\frac{\xi_{j}^{n}}{\Delta S}-\left(\varphi+c_{n}\right) \frac{w_{j}^{n}}{\Delta S}\right)^{2} \sin ^{2}(q \Delta S)\right] \\
\leq & 9 \Delta \tau^{2} C_{3}^{2}\left|\nu^{1}\right|^{2}, \Rightarrow\left|\vartheta^{1}\right| \leq 3 \Delta \tau C_{3}\left|\nu^{1}\right| \leq C_{3}(1+3 \Delta \tau)\left|\nu^{1}\right|
\end{aligned}
$$

Now, we suppose $\left|\vartheta^{k}\right| \leq C_{3}(1+3 \Delta \tau)^{k}\left|\nu^{1}\right|, k=2,3, \ldots, n-1$. Applying (4.10), Lemma 4.3 and (4.6) into (4.8), we prove $\left|\vartheta^{n}\right| \leq C_{3}(1+3 \Delta \tau)^{n}\left|\nu^{1}\right|$, where $C_{3}=\max \left\{C_{2}, C_{3}\right\}$,

$$
\begin{aligned}
\left|\vartheta^{n}\right| \leq & \left|\theta_{j}^{n} \varphi+i \frac{w_{j}^{n}}{\Delta S} \varphi \sin (q \Delta S)-\frac{\varphi}{3} \sin ^{2}\left(\frac{q \Delta S}{2}\right)\right| / \\
& +i\left(\frac{\xi_{j}^{n}}{\Delta S}-\left(\varphi+c_{n}\right) \frac{w_{j}^{n}}{\Delta S}\right) \sin (q \Delta S)\left|\times\left|\sum_{k=1}^{n-1}\left(\psi_{k+1}^{n}-\psi_{k}\right) \vartheta^{n-k}\right|\right. \\
& +\left|\nu^{n}\right| / \left\lvert\,\left(-\frac{4 \gamma_{j}^{n}}{\Delta S^{2}}+\frac{\varphi+c_{n}}{3}\right) \sin ^{2}\left(\frac{q \Delta S}{2}\right)-\left(\varphi+2 c_{n}\right) \theta_{j}^{n}\left(\frac{q \Delta S}{2}\right)-\left(\varphi+2 c_{n}\right) \theta_{j}^{n}\right. \\
& \left.+i\left(\frac{\xi_{j}^{n}}{\Delta S}-\left(\varphi+c_{n}\right) \frac{w_{j}^{n}}{\Delta S}\right) \sin (q \Delta S) \right\rvert\,
\end{aligned}
$$

$$
\begin{aligned}
& \leq\left|\sum_{k=1}^{n-1}\left(\psi_{k+1}-\psi_{k}\right) \vartheta^{n-k}\right|+3\left|\nu^{n}\right| \\
& \leq \sum_{k=1}^{n-1}\left(\psi_{k}-\psi_{k+1}\right)\left|\vartheta^{n-k}\right|+3 C_{2} \Delta \tau\left|\nu^{1}\right| \\
& \leq \sum_{k=1}^{n-1}\left(\psi_{k}-\psi_{k+1}\right) C_{3}(1+3 \Delta \tau)^{n-k}\left|\nu^{1}\right|+C_{2}(1+3 \Delta \tau)\left|\nu^{1}\right| \\
& \leq C_{3}(1+3 \Delta \tau)^{n-1}\left|\nu^{1}\right| \sum_{k=1}^{n-1}\left(\psi_{k}-\psi_{k+1}\right)+C_{2}(1+3 \Delta \tau)\left|\nu^{1}\right| \\
& \leq C_{3}(1+3 \Delta \tau)^{n-1}\left|\nu^{1}\right|\left(1-\psi_{n}\right)+C_{2}(1+3 \Delta \tau) \psi_{n}\left|\nu^{1}\right| \\
& \leq C_{3}(1+3 \Delta \tau)^{n}\left|\nu^{1}\right|\left(1-\psi_{n}+\psi_{n}\right)=C_{3}(1+3 \Delta \tau)^{n}\left|\nu^{1}\right| .
\end{aligned}
$$

This ends the proof.
Theorem 4.8. Assume that $V(S, \tau)$ is the exact solution of (2.5) and $\widetilde{V}(S, \tau)$ is the exact solution of (3.8), the compact difference scheme (3.8) is convergent, and the convergence order is $O\left(\Delta \tau^{2-\alpha}+\Delta S^{4}\right)$.
Proof. Consider Lemma 4.7, combine (4.7a), (4.7b) and (4.9)

$$
\left\|E^{n}\right\|_{2} \leq C_{3}(1+3 \Delta \tau)^{n}\left\|R^{1}\right\|_{2} \leq C_{1} C_{3} \sqrt{L} \exp (3 n \Delta \tau)\left(\Delta \tau^{2-\alpha}+\Delta S^{4}\right) .
$$

Since $n \Delta \tau \leq T$, we derive $\left\|E^{n}\right\|_{2} \leq C\left(\Delta \tau^{2-\alpha}+\Delta S^{4}\right)$, where

$$
C=C_{1} C_{3} \sqrt{L} \exp (3 n \Delta \tau),
$$

This finishes the proof.

## 5. Numerical examples

Now, we exhibit the accuracy of the introduced scheme with three examples for solving the time-fractional Black-Scholes equation under CEV model in which interest rate and dividend yield are as deterministic time-dependent parameters. We compute the time of running the program by using CPU times of MATLAB R2015a. The CPU times show the low volume of computation and the advantages of the introduced scheme. Furthermore, in all three examples, we will compare the generalized BlackScholes model with the time-fractional Black-Scholes equation under CEV model when $\alpha=1$ and $\beta=0$. We also discuss the Greek letters $\Delta, \Gamma$ and $\Theta$ using the figure.
Example 5.1. Consider the time-fractional Black-Scholes Eq. (2.5) with the parameters: $r(t)=0.1+0.05 e^{-t}, D(t)=0.03+0.001 e^{0.01 t}[50], K=50, \sigma_{0}=0.4, S_{0}=50$, $T=3, \beta=-0.5, \alpha=0.8, N=M=100$ and $S_{\max }=3 K$.

Example 5.2. Price the American call option model (2.5) with the parameters: $r(t)=0.075+0.05 t[29], D(t)=0.05, K=50, \sigma_{0}=0.4, S_{0}=50, T=3, \beta=-0.5$, $\alpha=0.8, N=M=100$ and $S_{\max }=3 K$.

Example 5.3. Obtain the American call option price for model (2.5) with the parameters: $r(t)=0.1+0.0005 t[39], D(t)=0.02, K=50, \sigma_{0}=0.4, S_{0}=50, T=3$, $\beta=-0.5, \alpha=0.8, N=M=100$ and $S_{\max }=3 K$.

TABLE 1. CPU time to determine option price in expiry date.

| N | M | Example 1 | Example 2 | Example 3 |
| :---: | :---: | :---: | :---: | :---: |
|  |  | CPU time | CPU time | CPU time |
| 64 | 64 | 0.281639 s | 0.286932 s | 0.292268 s |
| 128 | 128 | 1.242377 s | 1.221253 s | 1.203057 s |
| 256 | 264 | 4.881068 s | 4.815844 s | 4.887891 s |
| 512 | 512 | 23.071165 s | 22.823704 s | 23.102432 s |
| 1024 | 1024 | 128.594945 s | 130.967890 s | 130.980526 s |

TABLE 2. Option price for different $\alpha$ in expiry date.

|  |  |  | $\alpha=0.80$ | $\alpha=0.85$ | $\alpha=0.90$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\alpha=0.95$ |  |  |  |  |
| Example 1 | 30 | 5.1104 | 5.3491 | 5.5929 | 5.8421 |
|  | 60 | 22.4017 | 22.8049 | 23.2154 | 23.6336 |
|  | 90 | 49.0168 | 49.2749 | 49.5398 | 49.8124 |
|  | 120 | 80.8191 | 80.9912 | 81.1681 | 81.3502 |
| Example 2 | 30 | 5.0655 | 5.3655 | 5.6727 | 5.9873 |
|  | 60 | 21.5117 | 22.0903 | 22.6809 | 23.2838 |
|  | 90 | 46.7652 | 47.2432 | 47.7340 | 48.2382 |
|  | 120 | 76.4837 | 76.8334 | 77.1894 | 77.5520 |
| Example 3 | 30 | 5.1687 | 5.4266 | 5.6891 | 5.9565 |
|  | 60 | 22.4865 | 22.9615 | 23.4440 | 23.9345 |
|  | 90 | 48.8599 | 49.2118 | 49.5711 | 49.9387 |
|  | 120 | 79.2868 | 79.5290 | 79.7757 | 80.0272 |

Table 3. Option price for different $\beta$ in expiry date.

|  | $S$ | $\beta=-0.4$ | $\beta=-0.3$ | $\beta=-0.2$ | $\beta=-0.1$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Example 1 | 30 | 4.9234 | 4.7464 | 4.5783 | 4.4180 |
|  | 60 | 22.5520 | 22.7105 | 22.8762 | 23.0474 |
|  | 90 | 49.3856 | 49.7679 | 50.1598 | 50.5572 |
|  | 120 | 81.2265 | 81.6356 | 82.0426 | 82.4436 |
| Example 2 | 30 | 4.8568 | 4.6594 | 4.4723 | 4.2943 |
|  | 60 | 21.6257 | 21.7473 | 21.8753 | 22.0081 |
|  | 90 | 47.0949 | 47.4342 | 47.7789 | 48.1251 |
|  | 120 | 76.8457 | 77.2025 | 77.5505 | 77.8865 |
| Example 3 | 30 | 4.9648 | 4.7716 | 4.5878 | 4.4126 |
|  | 60 | 22.6000 | 22.7194 | 22.8439 | 22.9724 |
|  | 90 | 49.1333 | 49.4172 | 49.7085 | 50.0039 |
|  | 120 | 79.5577 | 79.8313 | 80.1045 | 80.3747 |

We compute CPU time of Examples 5.1-5.3 for different $N$ and $M$ in Table 1. This table shows that CPU time is almost 2 min 10 s for $N=M=1024$. We investigate the effect of each parameter time-fractional derivative order $(\alpha)$, elasticity factor $(\beta)$, and initial instantaneous volatility $\left(\sigma_{0}\right)$ on the long memory in Tables 2-4. Table 2 displays that option price is increasing for $\alpha=\{0.80,0.85,0.90,0.95\}$. Table 3 shows that option price is both decreasing and increasing for $\beta=\{-0.4,-0.3,-0.2,-0.1\}$. Table 4 is increasing for $\sigma_{0}=\{0.2,0.3,0.4,0.5\}$.

TABLE 4. Option price for different $\sigma_{0}$ in expiry date.

|  | $S$ | $\sigma_{0}=0.2$ | $\sigma_{0}=0.3$ | $\sigma_{0}=0.4$ | $\sigma_{0}=0.5$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Example 1 | 30 | 1.6013 | 3.2722 | 5.1104 | 7.0233 |
|  | 60 | 18.7947 | 20.2931 | 22.4017 | 24.8057 |
|  | 90 | 46.5690 | 47.5341 | 49.0168 | 50.8966 |
|  | 120 | 78.1671 | 79.5494 | 80.8191 | 82.1093 |
| Example 2 | 30 | 1.7016 | 3.3045 | 5.0655 | 6.9024 |
|  | 60 | 17.9890 | 19.4581 | 21.5117 | 23.8501 |
|  | 90 | 44.4315 | 45.3374 | 46.7652 | 48.5862 |
|  | 120 | 73.9721 | 75.2834 | 76.4837 | 77.6991 |
| Example 3 | 30 | 1.5689 | 3.2812 | 5.1687 | 7.1194 |
|  | 60 | 18.7584 | 20.3488 | 22.4865 | 24.8718 |
|  | 90 | 46.9493 | 47.6299 | 48.8599 | 50.5255 |
|  | 120 | 77.5483 | 78.4229 | 79.2868 | 80.2690 |

TABLE 5. Convergence rate for different $N$ when $K=40, \sigma_{0}=0.4$, $S_{0}=40, T=3, \beta=-1, \alpha=0.7$ and $S_{\max }=3 K$.

| $r(t)=0.1+0.05 e^{-t}, D(t)=0.03+0.001 e^{0.01 t}$ |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
| M | N | Error | Rate | CPU time |
| 1024 | 64 | $1.5287 \mathrm{e}-05$ | - | - |
|  | 128 | $9.9149 \mathrm{e}-07$ | 3.9466 | 551.254211 s |
|  | 256 | $6.4936 \mathrm{e}-08$ | 3.9325 | 561.611072 s |
|  | 512 | $4.0624 \mathrm{e}-09$ | 3.9986 | 597.201253 s |
|  | 1024 | $2.4475 \mathrm{e}-10$ | 4.0529 | 660.541270 s |
| $r(t)=0.075+0.05 t, D(t)=0.05$ |  |  |  |  |
| M | N | Error | Rate | CPU time |
| 1024 | 64 | $1.1963 \mathrm{e}-05$ | - | - |
|  | 128 | 7.7641e-07 | 3.9456 | 551.784606 s |
|  | 256 | $5.0150 \mathrm{e}-08$ | 3.9525 | 569.122115 s |
|  | 512 | $3.2328 \mathrm{e}-09$ | 3.9554 | 590.740887 s |
|  | 1024 | $1.8889 \mathrm{e}-10$ | 4.0972 | 656.945209 s |
| $r(t)=0.1+0.0005 t, D(t)=0.02$ |  |  |  |  |
| M | N | Error | Rate | CPU time |
| 1024 | 64 | $6.9990 \mathrm{e}-06$ | - | - |
|  | 128 | $2.8956 \mathrm{e}-07$ | 4.5952 | 540.760544 s |
|  | 256 | $2.0692 \mathrm{e}-08$ | 3.8067 | 550.976153 s |
|  | 512 | $1.3045 \mathrm{e}-09$ | 3.9875 | 578.041359 s |
|  | 1024 | $7.5417 \mathrm{e}-11$ | 4.1124 | 664.398942 s |

The American options have no closed-form solution. Therefore, to illustrate the fourth-order convergence rate in space numerically, we compare our solution with the approximated solution that $N$ and $M$ are large enough (see [5]). We define the discrete Maximum-norm error as

$$
e^{N, M}=\max _{\substack{0 \leq j \leq N \\ 0 \leq n \leq M}}\left\|\bar{V}_{j}^{n}-\widehat{V}_{j}^{n}\right\|_{\infty}=\max _{0 \leq n \leq M}\left(\max _{0 \leq j \leq N}\left|\bar{V}_{j}^{n}-\widehat{V}_{j}^{n}\right|\right)
$$

where $\bar{V}_{j}^{n}$ is the approximated solution for $N=2048$ and $M=1024$ and $\widehat{V}_{j}^{n}$ is the our solution. The convergence rate in space is obtained from following relation

$$
\text { Rate }=\log _{2}\left(\frac{e^{N, M}}{e^{2 N, M}}\right) .
$$

Table 6. Comparison of the American put option price by compact difference scheme with Zhou's result at $r=0.05, D=0, K=40, S_{0}=40$, $T=3, S_{\max }=200, N=800$ and $M=200$ in expiry date.

| $\beta=0$ | $\sigma_{0}=0.1$ |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  | $\sigma_{0}=0.2$ |  |  |  |
| $\alpha$ | LTM | FDM | Compact |  | LTM | FDM | Compact |
| 1.0 | 1.2189 | 1.2362 | 1.2308 |  | 3.4116 | 3.4792 | 3.4741 |
| 0.9 | 1.1771 | 1.1912 | 1.1830 |  | 3.2651 | 3.3157 | 3.3080 |
| 0.7 | 1.0959 | 1.1028 | 1.0879 |  | 2.9817 | 3.0071 | 2.9927 |
| 0.4 | 0.9778 | 0.9793 | 0.9347 |  | 2.5770 | 2.5829 | 2.5418 |
| 0.2 | 0.9005 | 0.9002 | 0.8094 |  | 2.3184 | 2.3191 | 2.2341 |
| $\beta=-1$ |  |  |  |  |  |  |  |
|  |  | $\sigma_{0}=0.1$ |  |  | $\sigma_{0}=0.2$ |  |  |
| $\alpha$ | LTM | FDM | Compact |  | LTM | FDM | Compact |
| 1.0 | 1.1877 | 1.2020 | 1.1977 |  | 3.3325 | 3.3834 | 3.3880 |
| 0.9 | 1.1485 | 1.1604 | 1.1528 |  | 3.1918 | 3.2297 | 3.2300 |
| 0.7 | 1.0722 | 1.0802 | 1.0641 |  | 2.9208 | 2.9400 | 2.9309 |
| 0.4 | 0.9609 | 0.9657 | 0.9202 |  | 2.5347 | 2.5397 | 2.5048 |
| 0.2 | 0.8877 | 0.8922 | 0.8017 |  | 2.2876 | 2.2898 | 2.2158 |

Table 7. Comparison of the American put option price by compact difference scheme with Pun's result at $\alpha=1, r=0.05, D=0, \sigma_{0}=0.4$, $\beta=-0.1, K=40, S_{0}=40, T=1$ and $S_{\max }=3 K$ in expiry date.

| Pun's result |  |  |  |  | $\begin{gathered} \alpha=1 \\ \mathrm{~N}=\mathrm{M}=240 \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| S | BS sol. | 1st-order | 2nd-order | RE | Compact |
| 20 | 19.7979 | 20 | 20 | 20 | 20 |
| 30 | 12.6119 | 11.1815 | 10.7256 | 10.7299 | 10.9832 |
| 40 | 7.98526 | 5.77404 | 5.36116 | 5.33151 | 5.4557 |
| 50 | 5.20209 | 2.84158 | 2.58372 | 2.54957 | 2.5448 |
| 60 | 3.38609 | 1.23602 | 1.18842 | 1.13693 | 1.1450 |

In Table 5, we list the error estimates and convergence rates of the introduced scheme for three $r(t)$ and $D(t)$. This table shows that the obtained convergence rates support Theorem 4.8. The CPU times in this table represent the sum of the run times associated with computing $e^{N, M}, e^{2 N, M}$ and convergence rate per run.

We can also use the introduced difference scheme to determine the American put option price. Hence, we present two comparisons of the introduced scheme with the results of [53] and [41] in Tables 6 and 7, respectively.


Figure 1. Option price of three examples in expiry date with their payoff.

Figure 1 shows the option price of three examples in the expiry date with their payoff. Figure 2 displays that time-fractional Black-Scholes equations under the CEV model equal to generalized Black-Scholes model in expiry date when $\alpha=1$ and $\beta=0$ for all examples. The generalized Black-Scholes model of this figure with an implicit difference scheme is described in Appendix A. Figures 3-5 illustrate the effect of parameters $\alpha, \beta$, and $\sigma_{0}$ on option price, respectively. Figures $6-8$ represent option price sensitivities relative to the parameters.

## 6. Conclusion

Due to the limitations of the Black-Scholes model, we need a model that is closer to market realities and show memory effect in financial pricing. In this work, we investigated American call option pricing based on the time-fractional Black-Scholes equation under the CEV model with time-dependent parameters of risk-free interest rate and dividend yield. We presented a compact difference scheme to price the American call option as numerically. We analyzed stability and convergence of the introduced difference scheme using Fourier analysis and showed that the introduced scheme has the fourth-order convergence rate in space. Numerical examples express that the time-fractional Black-Scholes equation under the CEV model coincides with its generalized Black-Scholes equation as $\alpha=1$ and $\beta=0$. Also, we observed which American option price is increasing with respect to the time-fractional order derivative $(\alpha)$ and initial instantaneous volatility $\left(\sigma_{0}\right)$, and the American option price is both


Figure 2. Comparison of generalized Black-Scholes model and timefractional Black-Scholes equation under CEV model in expiry date when $\alpha=1$ and $\beta=0$.
decreasing and increasing as elasticity factor $(\beta)$ is increasing. Moreover, we discussed the American option price sensitivities relative to the underlying asset and time to expire. As a suggestion, we can use the introduced difference scheme of this paper to price other options.

## Appendix A. Generalized Black-Scholes model its difference scheme

The generalized Black-Scholes model for pricing American call option is following form

$$
\begin{align*}
& \frac{\partial C(S, t)}{\partial t}+\frac{1}{2} \sigma^{2} S^{2} \frac{\partial^{2} C(S, t)}{\partial S^{2}}+(r(t)-D(t)) S \frac{\partial C(S, t)}{\partial S} \\
& \quad-r(t) C(S, t)=0  \tag{A.1}\\
& C(S, T)=\max (S-K, 0), \\
& C(0, t)=0 \\
& C\left(S_{f}(t), t\right)=S_{f}(t)-K, \\
& \frac{\partial C\left(S_{f}(t), t\right)}{\partial S}=1
\end{align*}
$$



Figure 3. Option price based on different $\alpha$ in expiry date.


Figure 4. Option price based on different $\beta$ in expiry date.


Figure 5. Option price based on different $\sigma_{0}$ in expiry date.


Figure 6. American options in the Greeks $\Delta$.


Figure 7. American options in the Greeks $\Gamma$.
where $\sigma$ is volatility of underlying asset. Using approximation derivatives,

$$
\begin{aligned}
& \frac{\partial C\left(S_{j}, t_{n}\right)}{\partial t}=\frac{C\left(S_{j}, t_{n+1}\right)-C\left(S_{j}, t_{n}\right)}{\Delta t}+O(\Delta t) \\
& \frac{\partial^{2} C\left(S_{j}, t_{n}\right)}{\partial S^{2}}=\frac{C\left(S_{j-1}, t_{n}\right)-2 C\left(S_{j}, t_{n}\right)+C\left(S_{j+1}, t_{n}\right)}{\Delta S^{2}}+O\left(\Delta S^{2}\right), \\
& \frac{\partial C\left(S_{j}, t_{n}\right)}{\partial S}=\frac{C\left(S_{j+1}, t_{n}\right)-C\left(S_{j-1}, t_{n}\right)}{2 \Delta S}+O\left(\Delta S^{2}\right)
\end{aligned}
$$

we apply following implicit difference scheme on (A.1) as

$$
\begin{aligned}
& \frac{C_{j}^{n+1}-C_{j}^{n}}{\Delta t}+\frac{1}{2} \sigma^{2} S_{j}^{2}\left[\frac{C_{j-1}^{n}-2 C_{j}^{n}+C_{j+1}^{n}}{2 \Delta S^{2}}+\frac{C_{j-1}^{n+1}-2 C_{j}^{n+1}+C_{j+1}^{n+1}}{2 \Delta S^{2}}\right] \\
& \quad+\left[r\left(t_{n}\right)-D\left(t_{n}\right)\right] S_{j}\left[\frac{C_{j+1}^{n}-C_{j-1}^{n}}{4 \Delta S}+\frac{C_{j+1}^{n+1}-C_{j-1}^{n+1}}{4 \Delta S}\right] \\
& \quad-r\left(t_{n}\right)\left[\frac{C_{j}^{n}+C_{j}^{n+1}}{2}\right]=0 .
\end{aligned}
$$

The matrix form above scheme can be written as

$$
A^{n} C^{n+1}=B^{n} C^{n}-F^{n+1}, \quad n \geq 0,
$$



Figure 8. American options in the Greeks $\Theta$.
where

$$
\begin{aligned}
& A^{n}=\left[\begin{array}{ccccc}
b_{1}^{n} & e_{2}^{n} & & & 0 \\
a_{1}^{n} & b_{2}^{n} & e_{3}^{n} & & \\
& \ddots & \ddots & \ddots & \\
& & & & e_{N-1}^{n} \\
0 & & & a_{N-2}^{n} & b_{N-1}^{n}
\end{array}\right] \\
& B^{n}=\left[\begin{array}{ccccc}
d_{1}^{n} & -e_{2}^{n} & & & 0 \\
-a_{1}^{n} & d_{2}^{n} & -e_{3}^{n} & & \\
& \ddots & \ddots & \ddots & \\
& & & & -e_{N-2}^{n} \\
0 & & & d_{N-1}^{n}
\end{array}\right] \\
& C^{n}=\left[C_{1}^{n}, C_{2}^{n}, \ldots, C_{N-1}^{n}\right]^{t}, \\
& F^{n+1}=\left[a_{0}^{n}\left(C_{0}^{n}+C_{0}^{n+1}\right), 0, \ldots, 0, e_{N}^{n}\left(C_{N}^{n}+C_{N}^{n+1}\right)\right]^{t} .
\end{aligned}
$$

and

$$
\left\{\begin{array}{l}
a_{j-1}^{n}=\frac{\sigma^{2}}{4} \frac{\Delta t}{\Delta S^{2}} S_{j}^{2}-\frac{\Delta t}{4 \Delta S} S_{j}\left[r\left(t_{n}\right)-D\left(t_{n}\right)\right] \\
b_{j}^{n}=-\frac{\sigma^{2}}{2} \frac{\Delta t}{\Delta S^{2}} S_{j}^{2}-\frac{\Delta t}{2} r\left(t_{n}\right)+1, \\
e_{j+1}^{n}=\frac{\sigma^{2}}{4} \frac{\Delta t}{\Delta S^{2}} S_{j}^{2}+\frac{\Delta t}{4 \Delta S} S_{j}\left[r\left(t_{n}\right)-D\left(t_{n}\right)\right], \\
d_{j}^{n}=\frac{\sigma^{2}}{2} \frac{\Delta t}{\Delta S^{2}} S_{j}^{2}+\frac{\Delta t}{2} r\left(t_{n}\right)+1,
\end{array}\right.
$$

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